

UARS Microwave Limb Sounder HNO₃ observations: Implications for Antarctic polar stratospheric clouds

M. L. Santee,¹ A. Tabazadeh,² G. L. Manney,¹ R. J. Salawitch,¹ L. Froidevaux,¹
W. G. Read,¹ and J. W. Waters¹

Abstract. We present Microwave Limb Sounder (MLS) measurements of gas-phase HNO₃ obtained at the beginning of five southern hemisphere winters: 1992–1996. To investigate the composition of the polar stratospheric clouds (PSCs) that formed in early winter each year, the observed evolution of HNO₃ at 465 K is compared against that predicted using nitric acid trihydrate (NAT), nitric acid dihydrate (NAD), and liquid ternary solution models of PSC formation and correlated with temperature histories from three-dimensional back trajectory calculations. The MLS HNO₃ observations suggest that the initial composition of PSCs depends on the physical state of the background sulfate aerosols. If the preexisting aerosols are liquid, then the formation of ternary solutions is initiated as the temperature drops below about 192 K, followed by a gradual conversion to NAD after exposure to low temperatures for several days. HNO₃ uptake into ternary solutions occurs at higher temperatures, and the conversion to NAD is delayed, under conditions of enhanced aerosol loading from the Mount Pinatubo eruption. If a majority of the background aerosols are frozen, the growth of ternary solutions is inhibited, but formation of a metastable, water-rich, HNO₃-containing solid phase characterized by a relatively high HNO₃ vapor pressure (type Ic PSC) may occur. In general, MLS HNO₃ measurements obtained during early southern winter indicate a strong correspondence between the area of gas-phase HNO₃ loss and the area of temperatures below 192 K but only a weak correspondence between the area of gas-phase HNO₃ loss and the area of temperatures below 195 K, the value commonly assumed as the threshold for PSC formation. Although temperatures were low enough to maintain NAT PSCs, the MLS data show that they were not forming, at least not over spatial scales comparable to or larger than the $\sim 400 \times 100 \times 5$ km MLS field of view.

1. Introduction

It is well established that the chlorine-catalyzed destruction of ozone in the polar lower stratosphere during winter and spring is initiated through heterogeneous chemical reactions on the surfaces of polar stratospheric cloud (PSC) particles [e.g., Solomon, 1990]. Accurate theoretical predictions of ozone loss therefore require knowledge of the phase and composition of the PSC particles. For instance, recent work [Ravishankara and Hanson, 1996; Borrmann *et al.*, 1997] has shown that liquid particles may activate chlorine more efficiently than frozen particles. However, despite a considerable number of field, laboratory, and modeling studies, the exact PSC formation mechanism remains uncertain.

There are two main classifications of PSCs: type I and type II. Type II PSCs are composed of water ice and typically form a few degrees below the equilibrium condensation

point of ice (nominally 188 K for lower stratospheric conditions) because supersaturations with respect to ice exceeding unity are required to overcome the free-energy barrier to nucleation [e.g., Toon *et al.*, 1989]. In contrast, type I PSCs are observed to form at temperatures several degrees above the water ice frost point [e.g., McCormick *et al.*, 1982]. Toon *et al.* [1986] and Crutzen and Arnold [1986] suggested that type I PSCs contain HNO₃, and the principal phase was usually assumed to be nitric acid trihydrate (NAT, HNO₃·3H₂O), which condenses at ~ 195 K and which is the most stable form under stratospheric conditions [Hanson and Mauersberger, 1988]. Early field observations [e.g., Fahney *et al.*, 1989; Poeschel *et al.*, 1989] confirmed that HNO₃ is a major component of type I PSC particles and that the temperature at which they first appear is roughly consistent with formation of a trihydrate phase. For years the canonical model of PSC formation was predicated on a NAT composition for the type I particles [e.g., Solomon, 1990].

However, there is a growing body of evidence that not all type I PSC particles are composed of NAT. In the Arctic, large NAT supersaturations have been observed before the onset of significant particle growth [e.g., Dye *et al.*, 1992], and the amount of HNO₃ incorporated into the particles has been found to be significantly less than that predicted for a NAT composition [Kawa *et al.*, 1992]. Two

¹Jet Propulsion Laboratory, California Institute of Technology, Pasadena.

²NASA Ames Research Center, Moffett Field, California.

subclasses of type I PSCs have been recognized based on lidar measurements of aerosol optical properties [Browell *et al.*, 1990; Toon *et al.*, 1990a]: type Ia, relatively large, nonspherical particles, and type Ib, smaller, nearly spherical particles. Laboratory measurements [Zhang *et al.*, 1993a; Molina *et al.*, 1993] have shown that HNO₃ is highly soluble in aqueous H₂SO₄ at the low temperatures typical of polar winter, and a supercooled ternary (H₂SO₄/HNO₃/H₂O) solution composition has been proposed for type Ib particles [Tabazadeh *et al.*, 1994a]. Previously obtained field observations have now been reinterpreted as being more consistent with a ternary than with a NAT composition for both the Arctic [Drdla *et al.*, 1994; Carslaw *et al.*, 1994; Tabazadeh *et al.*, 1994a] and the Antarctic [Toon and Tolbert, 1995]. Ground-based lidar [Beyerle *et al.*, 1997] and balloon-borne backscatter [Larsen *et al.*, 1997] measurements have provided further evidence for liquid ternary solution PSCs in the Arctic. In addition, analyses of total reactive nitrogen and aerosol particle volume measured during the 1994 Airborne Southern Hemisphere Ozone Experiment (ASHOE) have indicated that a type I PSC near the edge of the Antarctic vortex was composed primarily of ternary solution droplets [Dye *et al.*, 1996; Del Negro *et al.*, 1997]. Trends in ground-based measurements of stratospheric HNO₃ over the south pole have also suggested the formation of ternary solution PSCs [de Zafra *et al.*, 1997].

Although formation of liquid type Ib PSCs is now thought to be understood, the mechanisms by which crystalline type Ia PSCs form in the stratosphere remain unresolved [Tolbert, 1994]. Previously, it was assumed that frozen sulfuric acid aerosols provide nucleation sites for the formation of type Ia PSCs [e.g., Turco *et al.*, 1989; Toon *et al.*, 1989]. However, laboratory experiments on sulfuric acid films [Middlebrook *et al.*, 1993] have indicated that temperatures near the ice frost point are necessary for the stratospheric sulfate aerosol to freeze, forming sulfuric acid tetrahydrate (SAT); recent experiments on binary sulfate aerosol particles [Imre *et al.*, 1997; Carleton *et al.*, 1997; Clapp *et al.*, 1997] have confirmed that stratospheric sulfate aerosols most likely remain liquid to the ice frost point or below. Laboratory observations of ternary bulk solutions [Koop *et al.*, 1995] and freely floating ternary droplets [Anthony *et al.*, 1997] have also found that supercooled ternary solutions do not freeze for temperatures above the ice frost point. Other laboratory studies [Molina *et al.*, 1993; Beyer *et al.*, 1994; Iraci *et al.*, 1995] have shown that NAT crystallization directly from the vapor phase onto SAT is unlikely; these studies suggest that NAT nucleates out of liquid ternary solutions, with subsequent growth of the crystals by condensation of HNO₃ and H₂O vapors. However, model calculations [Tabazadeh and Toon, 1996] indicate that the highest possible HNO₃ concentration in a ternary droplet is insufficient to allow the direct crystallization of NAT. Several candidates have been proposed for an initial metastable phase for condensed HNO₃ in PSCs, with eventual conversion to NAT. For example, laboratory results of Marti and Mauersberger [1993] suggest a dilute HNO₃/H₂O solid phase (with a water to nitric acid ratio of about 5 or 6 to 1)

as a transient PSC composition. Recent laboratory experiments by Fox *et al.* [1995] also suggest that NAT nucleation from a ternary solution requires the formation of a precursor phase, which they infer to be a mixed H₂SO₄/HNO₃ hydrate (MixH, H₂SO₄·HNO₃·5H₂O). In cooling a ternary solution under stratospheric conditions, Fox *et al.* [1995] never observed NAT nucleation in the solution without prior formation of MixH, which then transformed into a metastable nitric acid dihydrate (NAD, HNO₃·2H₂O) phase within a few hours. NAT was always the last phase to form, with NAD persisting for days, in agreement with their earlier laboratory experiments on PSCs [Worsnop *et al.*, 1993] showing NAD to be only slightly less stable than NAT.

Alternatively, laboratory studies by Koop *et al.* [1995] show that ternary bulk solutions remain supercooled to temperatures below the ice frost point. They contend that only water ice is a suitable nucleation site for H₂SO₄ and HNO₃ hydrates, and thus stratospheric aerosol droplets must cool below the frost point before crystallization of either HNO₃ or H₂SO₄ in solution is possible. Meilinger *et al.* [1995] suggest that the rapid cooling rates induced by orographically forced lee waves lead to nonequilibrium conditions that facilitate freezing of the smallest ternary droplets into NAT, since their composition is closest to that of trihydrate. However, recent freezing experiments on HNO₃/H₂O aerosols indicate that the HNO₃ concentration in solution must be larger than 1:2.5 (HNO₃:H₂O) in order for HNO₃ to crystallize [Disselkamp *et al.*, 1996]. This result argues against the Meilinger *et al.* [1995] proposal of crystallization of aerosols containing less than a 1:3 (HNO₃:H₂O) ratio. In addition, Tabazadeh *et al.* [1996] use aircraft observations of PSCs and a statistical analysis in which both background and lee wave temperature fluctuations are superimposed on synoptic-scale temperature histories to demonstrate that neither ice crystal nucleation nor lee wave encounters are correlated with the occurrence of type Ia PSCs. Murphy and Gary [1995] and Tabazadeh *et al.* [1996] argue that small but rapid temperature fluctuations (especially frequent in the Arctic) promote the initial formation of large numbers of small metastable particles with subsequent mass transfer to the larger, more stable NAT particles. In fact, Tabazadeh *et al.* [1996] conclude that mass transfer, and not the mesoscale cooling rate, controls the type Ia PSC formation process; once synoptic temperatures drop below about 192 K, the only condition necessary for type Ia formation is that temperatures remain below the NAT frost point for longer than ~1 day. Larsen *et al.* [1997] reach a similar conclusion based on balloon measurements of Arctic PSCs, finding that particle properties are most strongly influenced by the synoptic temperature history even for measurements obtained in regions where lee waves are common.

Building on the laboratory studies [Worsnop *et al.*, 1993; Marti and Mauersberger, 1993] indicating the initial presence of higher hydrates of HNO₃, Tabazadeh *et al.* [1995] and Tabazadeh and Toon [1996] have postulated a third subclass of type I PSCs, type Ic, composed of a water-rich metastable HNO₃/H₂O solid phase that transforms into NAT or NAD over time. In their scenario, the thermal history

of the air mass, and thus the physical state of the background sulfate aerosols, determines the phase of the PSCs: Upon cooling, if the H₂SO₄ aerosols are liquid, then type Ib ternary solutions form; if they are frozen, then type Ic particles initially form. Assuming that the ice frost point is not reached, if the cooling is then followed by a warming to about 196–198 K, some supercooled H₂SO₄/H₂O aerosols may freeze (as observed in the laboratory by *Iraci et al.* [1994]) and remain frozen until the air mass experiences temperatures high enough to melt SAT (about 210–220 K [Middlebrook *et al.*, 1993; Zhang *et al.*, 1993b]). With this mechanism, cooling/warming cycles eventually lead to the formation of type Ia PSCs. This scenario is consistent with aircraft observations of PSCs in both the Arctic and the Antarctic [Tabazadeh *et al.*, 1995, 1996] and is supported by analyses of balloon-borne backscatter measurements of PSCs from several Arctic winters [Larsen *et al.*, 1996, 1997]. Ground-based Arctic lidar measurements by Shibata *et al.* [1997] have also indicated the presence of large water-rich particles that are not compatible with the standard type Ia/b PSC classifications.

In this study we investigate PSC formation using measurements of gas-phase HNO₃ made by the Microwave Limb Sounder (MLS) aboard the Upper Atmosphere Research Satellite (UARS). Although these satellite measurements lack the resolution and precision of in situ measurements, they have the important advantage of providing daily hemispheric coverage over a range of pressures and over the course of several years. A description of the MLS HNO₃ measurements and their quality and spatial resolution is provided in section 2, and their suitability for studies of PSC composition is discussed in section 3. In section 4 we examine in detail the evolution of gas-phase HNO₃ during the 1994 early southern winter period, and we contrast this behavior with that during similar periods in 1992, 1993, 1995, and 1996. We compare the observed HNO₃ values in each year with those predicted by several proposed PSC composition models. Specifically, we discuss the MLS HNO₃ data in relation to type Ia PSCs based on the work of *Hanson and Mauersberger* [1988] and *Worsnop et al.* [1993], type Ib PSCs based on the work of *Tabazadeh et al.* [1994b], and type Ic PSCs based on the work of *Tabazadeh and Toon* [1996]; no new PSC subclasses are proposed here. Any comparison of MLS data with model results necessarily entails several assumptions; the observational and/or modeling studies underlying our assumptions, and the sensitivity of our results to them, are also discussed in section 4. Finally, in section 5 we use back trajectory calculations to explore the correlation between temperature history and PSC composition as inferred from MLS measurements of gas-phase HNO₃.

We find that the MLS HNO₃ observations indicate variations in PSC composition with time and meteorological conditions. Although our results are strongly suggestive, it must be borne in mind that the coarse spatial resolution and precision of the data, together with the necessity of estimating unmeasured quantities, limit our ability to definitively identify PSC compositions.

2. Measurement Description

2.1. General MLS Information

MLS has been acquiring millimeter-wavelength emission measurements of the stratosphere in both hemispheres since late September 1991. Measurements are made as the instrument field of view (FOV) is vertically scanned through the atmospheric limb in a plane perpendicular to the UARS velocity. The microwave limb sounding technique and the MLS instrument are described in detail by *Waters* [1993] and *Barath et al.* [1993], respectively. Advantages of this technique include the measurement of thermal emission, so that observations can be obtained both day and night, and the use of long wavelengths, so that the data are not degraded by the presence of PSCs or other stratospheric aerosols. Validation of the MLS instrument calibration and the primary MLS measurements is described in a special issue of *Journal of Geophysical Research* (101 (D6), 9,539–10,476) dedicated to the evaluation of the UARS data.

The MLS pointing geometry together with the inclination of the UARS orbit leads to a measurement latitudinal coverage extending from 80° on one side of the equator to 34° on the other. The UARS orbit plane precesses in such a way that all local solar times are sampled in ~36 days (a “UARS month”), getting ~20 min earlier each day at a given latitude. To keep MLS (and other instruments) on the shaded side of the spacecraft, a 180° yaw maneuver is performed at the end of every UARS month. Thus 10 times per year MLS alternates between viewing northern and southern high latitudes, with the first day of a particular UARS month occurring ~5 days earlier each year.

2.2. MLS HNO₃ Data Quality

Although HNO₃ was not a primary MLS measurement objective, a significant HNO₃ feature situated just outside the spectral region used to measure ozone imposes a weak slope through the 205-GHz ozone band that can be used to retrieve profiles of gas-phase HNO₃. MLS HNO₃ measurements (from precursory algorithms) were first presented by *Santee et al.* [1995]. With the release of MLS version 4 retrieval algorithms, HNO₃ is now a standard product available for every day of the mission on which MLS made measurements. Version 4 MLS HNO₃ data for the 1995–1996 and 1996–1997 northern winters have been presented by *Santee et al.* [1996] and *Santee et al.* [1997], respectively.

Preliminary validation studies indicate that the MLS HNO₃ data are scientifically useful on the 100-, 46-, and 22-hPa retrieval surfaces, where the estimated single-profile precisions are approximately 2.0, 3.0, and 4.5 ppbv, respectively. The precisions, estimated from several months of data, were based on the observed variability in a narrow latitude band centered around the equator that was selected to minimize the effects of natural atmospheric variability. The true precisions may be slightly better than these estimates since the actual atmospheric variation is not completely negligible. These empirical precisions are generally consistent with uncertainties derived theoretically by propagating

the measurement noise through the retrieval algorithm. Although the computed uncertainties consist of both random and systematic components, the random (noise) components usually dominate for HNO₃, which has a relatively weak signal (compared with ozone, for example). The noise contribution to the uncertainty can be reduced by averaging together individual measurements.

These calculated uncertainty values are included in the daily MLS Level 3A HNO₃ data files (the Level 3A files for each geophysical parameter consist of retrieved values on a vertical grid common to all UARS instruments producing atmospheric profiles). The retrieval algorithm, which is based on sequential estimation, employs an a priori estimate from a month- and latitude-dependent climatology developed by the UARS science team. Under conditions of poor measurement sensitivity when the contribution to the retrieved values from the a priori estimate exceeds 25%, the uncertainties are flagged with a negative sign. Quality control is imposed on the MLS HNO₃ data by discarding retrievals having negative uncertainty values. In addition, the Level 3P parameter file associated with each Level 3A file, which contains other diagnostics identifying bad data records, is examined, and only profiles for which the flag `MMAF_STAT` = "G," "T," or "t" and the flag `QUALITY_O3_205` = 4 (since HNO₃ is retrieved in the 205-GHz ozone band) are used. However, these quality control measures are not always sufficient to filter out large spikes in the retrieved HNO₃ mixing ratios; occasional anomalous retrievals are eliminated on an individual basis. More detailed information about the quality of the MLS data is available from the NASA Goddard Space Flight Center Distributed Access Archive Center (DAAC) as well as the MLS web page (<http://mils.jpl.nasa.gov>).

Initial comparisons (over a limited data set) with collocated UARS Cryogenic Limb Array Etalon Spectrometer (CLAES) HNO₃ observations [Kumer *et al.*, 1996] show that the MLS HNO₃ agrees well at 100 hPa but is usually 0–2 ppbv lower at 46 hPa and 0–4 ppbv higher at 22 hPa (particularly in the polar regions). However, even where biases between the two data sets exist, there is good correspondence in the morphology of the CLAES and MLS HNO₃ fields. Preliminary comparisons with other correlative data sets also indicate that MLS overestimates HNO₃ abundances near the profile peak (around 25 km). Further validation of the MLS HNO₃ data set is in progress.

2.3. Spatial Resolution of MLS HNO₃ Data

For optically thin cases such as the HNO₃ measurement, the relative contribution to the measured limb emission (the weighting function) along the observation path has a Gaussian distribution, the width of which sets the horizontal resolution along the line of sight to ~400 km [Waters, 1993]. The UARS orbital motion during the portion of the limb scan in which HNO₃ is measured smears the measurements over ~100 km in a direction perpendicular to the MLS line of sight. The FOV vertical extent at the tangent point is ~3 km, but version 4 data are produced on a vertical grid with points spaced three per decade in the logarithm of atmospheric pressure, yielding a vertical resolution of ~5 km.

Thus the retrieved MLS HNO₃ values at a particular geographic location and atmospheric level essentially represent averages over a ~400 × 100 × 5 km volume of air.

3. Using MLS Data to Study PSCs

An assumption inherent in using MLS observations of gas-phase HNO₃ depletion to infer PSC composition is that the instrument FOV is filled by a single PSC type. Therefore the suitability of the MLS data for these kinds of studies depends on the scale of the spatial inhomogeneity in the clouds. The prevalence of PSCs inside the winter polar vortices was first recognized in satellite solar occultation observations of aerosol extinction from the Stratospheric Aerosol Measurement (SAM) II experiment [McCormick *et al.*, 1982]. In particular, these observations show that Antarctic PSCs typically form at 20 km in the late May/early June time period and rapidly develop into extensive regions of enhanced extinction coincident with areas of low temperature [McCormick *et al.*, 1982, 1989; Poole and Pitts, 1994]. However, the extent of a typical cloud bank cannot be determined from SAM II data alone since the measurement coverage on a given day is restricted to a fixed latitude that is generally near the edge of the vortex in winter (away from the coldest regions in the vortex interior) and since the instrument field of view is ~200 km [McCormick *et al.*, 1989]. UARS CLAES measurements of aerosol extinction have also revealed widespread PSC activity [Mergenthaler *et al.*, 1993; Roche *et al.*, 1994; Mergenthaler *et al.*, 1997], but, again, these data have limited horizontal resolution (~500 km [Roche *et al.*, 1993]).

The spatial scale of a "typical" type I PSC is probably best estimated from airborne lidar data. Few such data have been reported for the Antarctic; although Browell *et al.* [1988] observed many PSCs of large vertical and horizontal extent (>5 km in altitude and >5°–10° in latitude), these data were collected on flights between late August and late September. Possibly of more relevance to PSC formation in the Antarctic early winter time period are the lidar data from several Arctic aircraft missions. PSCs were observed continuously over the entire latitude range from 76°N to 90°N in multiple layers between 19 km and 23 km during a flight on January 24, 1984 [McCormick *et al.*, 1985; Kent *et al.*, 1986]; although temperatures were sufficiently low that these may have been primarily water ice clouds. Two different sets of lidar data were obtained during the Airborne Arctic Stratospheric Expedition (AASE) of the winter of 1988–1989. McCormick *et al.* [1990] report lidar measurements from flights on January 24 and 31, 1989, that indicate type I PSCs spanning about 15° in latitude and 5 km in altitude. These measurements do not distinguish between type Ia and Ib PSCs, but they do reveal localized areas of type II PSCs embedded in the type I cloud. Data from a multiwavelength lidar system have been used [Browell *et al.*, 1990; Toon *et al.*, 1990a] to identify several type Ia PSCs extending continuously over 5° in latitude and 5 km in altitude during January 1989. On occasion, type Ib clouds were observed adjacent to type Ia clouds; although smaller in scale, the type Ib clouds were

still characterized by a depth greater than 2 km and a horizontal extent greater than 200 km [Browell *et al.*, 1990]. More isolated regions of type II PSCs were also observed. During another campaign, the European Arctic Stratospheric Ozone Experiment (EASOE), an airborne lidar detected a pervasive type I PSC on December 11, 1991, that extended over 600 km in the horizontal and 5 km in the vertical [Godin *et al.*, 1994]. Using the lidar measurements of Browell *et al.* [1990], Tabazadeh *et al.* [1996] found that at locations where type Ia PSCs were observed, synoptic temperatures had been below the NAT frost point for at least a day, and they concluded that synoptic temperature histories could be used to predict the occurrence of type Ia PSCs. If the main criterion for type Ia PSC formation is synoptic in nature, then type Ia PSCs themselves would be expected to exist over synoptic scales. Indeed, a new analysis of the Browell *et al.* [1990] lidar observations (O. B. Toon, personal communication, 1997) indicates the presence of continuous type Ia PSCs over vast areas from January 6 to February 2, 1989, during AASE.

PSCs have also been detected through in situ measurements of aerosol particle size and volume. Large-scale PSCs, greater than 1000 km in horizontal extent at ~ 20 km, have been observed in AASE particle data [Dye *et al.*, 1990, 1992; Kawa *et al.*, 1992]. The data for one of these PSC encounters (January 24, 1989) have been found to be consistent with a type Ib cloud [Drdla *et al.*, 1994; Carslaw *et al.*, 1994; Tabazadeh *et al.*, 1994a]. Similarly, a flight on July 28, 1994, during the ASHOE campaign intercepted a 500-km-long portion of a PSC composed of ternary droplets near the vortex edge at ~ 20 km [Dye *et al.*, 1996; Del Negro *et al.*, 1997].

Thus evidence accumulated from remote and in situ measurements indicates that both type Ia (crystalline) and type Ib (liquid) PSCs routinely extend continuously over spatial scales comparable to or larger than the MLS FOV. Although, in principle, the change in gas-phase HNO₃ due to a very patchy NAT PSC could, when averaged over the MLS FOV, appear instead to indicate a cloud of ternary droplets, the weight of evidence suggests that this scenario is unlikely (although patches of type II PSCs smaller than can be resolved by MLS are likely to be embedded in large type I clouds, particularly in the Antarctic). For the purposes of this investigation we have assumed that the PSCs are uniform over the MLS FOV.

Another complicating factor in using the MLS HNO₃ measurements to investigate PSC composition, in addition to the limited precision and spatial resolution of the data, is that some critical quantities are not measured and must be estimated. Interpretation of the data is therefore predicated on a number of assumptions (which are discussed in detail in section 4). For these reasons, our results regarding PSC composition, while suggestive, are not conclusive.

4. Observations and Model Results

To date, MLS has observed six southern hemisphere winters: 1992–1997. However, problems with the UARS so-

lar array caused the MLS instrument to be turned off during much of the early winter observing periods in both 1992 and 1995. In June 1995, UARS began operating in an instrument power sharing mode whereby MLS is periodically turned off; also at that time an MLS operational schedule consisting of (typically) 2 days of scanning followed by 1 day of “rest” was adopted to conserve the lifetime of the antenna scanning mechanism. The limited data coverage and the gaps in the data for these years are evident in Figure 1, which is discussed in more detail below. By the time of the 1997 early southern winter, the UARS power situation had degraded to the point that MLS measurements were collected on only 6 days during May and June; these data are not considered further here. Although they are presented briefly below, because of the short timescales involved, the sporadic data of 1992, 1995, and 1996 are not so well suited for a detailed study of PSC formation processes. In both 1993 and 1994, MLS obtained a complete set of daily limb scans during the early southern winter observing period. In this study we have chosen to examine the 1994 data more extensively for two reasons: the later onset of low temperatures and the earlier start of the observing period in 1994 (see Figure 1), capturing the initial stages of PSC formation, and the existence of data from the ASHOE campaign and the Polar Ozone and Aerosol Measurement (POAM) II instrument, providing supplementary information about environmental conditions in the polar vortex.

4.1. The 1994 Early Southern Winter

The behavior of MLS HNO₃ on selected days during the 1994 early southern winter south viewing period (May 24 to July 3) is shown in Plate 1. The MLS data are gridded by binning and averaging 24 hours of data and are then vertically interpolated to the 465-K potential temperature (θ) surface (ranging from ~ 30 to 60 hPa for the low temperatures inside the polar vortex) using United Kingdom Meteorological Office (UKMO) temperatures [Swinbank and O'Neill, 1994]. Contours of temperature and potential vorticity (PV) derived from the UKMO analyses are also shown. During southern late autumn and early winter (June–July), the lower stratospheric vortex is characterized by strong diabatic descent but very little latitudinal mixing across its boundary [Manney *et al.*, 1994]. In the absence of other effects, these conditions lead to large HNO₃ mixing ratios inside the vortex like those seen at the beginning of the observing period. The decreases observed in the HNO₃ abundances during winter are unlikely to be caused by dynamical processes and are attributed here to PSC activity. POAM II aerosol extinction measurements confirm the existence of PSCs during this time period [Fromm *et al.*, 1997].

Although 465-K temperatures have been below 195 K over a sizeable area since mid-May (see Figure 1), gas-phase HNO₃ values remain high in most regions of the vortex through the first week in June. A small pocket of low HNO₃ appears in the map on June 7 and intensifies as temperatures drop below 188 K the next day (not shown). Loss of gas-phase HNO₃ continues throughout the observing period, so that by the beginning of July very low HNO₃ concentrations

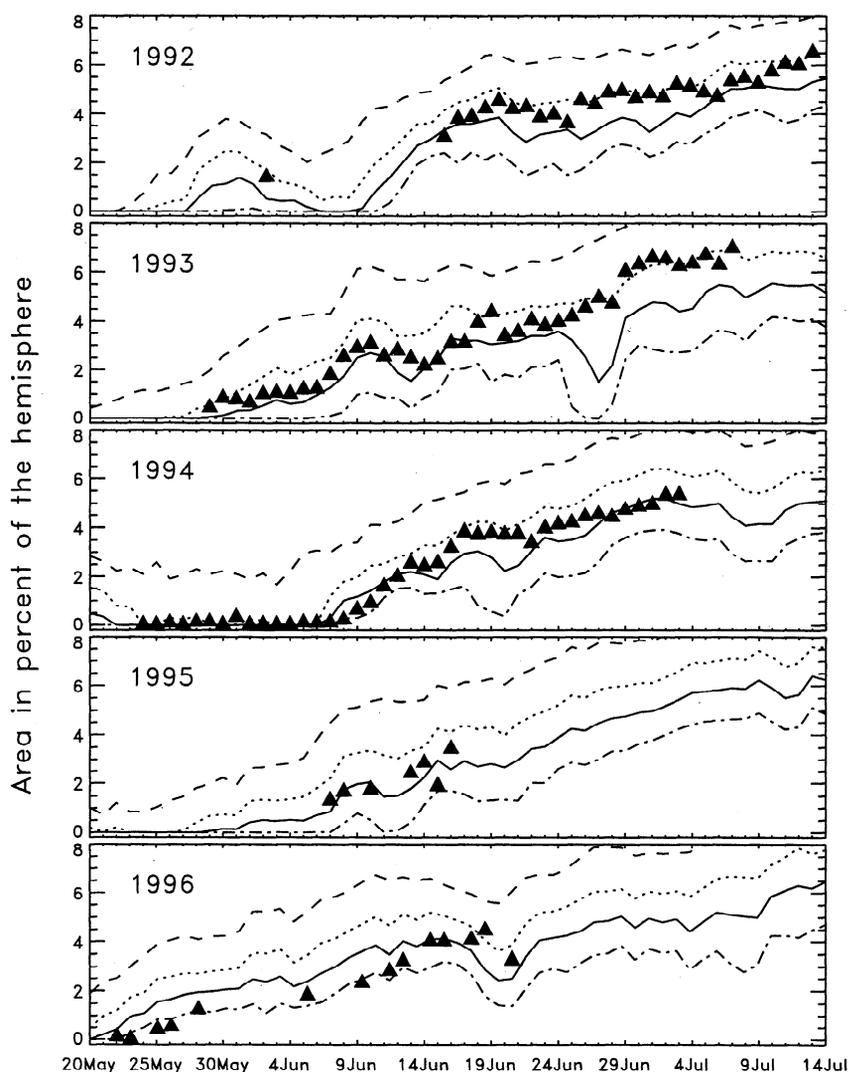


Figure 1. The area poleward of 60°S on the 465-K potential temperature surface (expressed as the percent of the hemisphere) as a function of time within which Microwave Limb Sounder (MLS) HNO₃ mixing ratios are below 4 ppbv (solid triangles) and United Kingdom Meteorological Office (UKMO) temperatures are below 188 K (dash-dotted lines), 190 K (solid lines), 192 K (dotted lines), and 195 K (dashed lines). Data for all five southern hemisphere early-winter south looking periods are shown. The solid triangles denote the days on which MLS performed full vertical scans in each year; data gaps in 1992, 1995, and 1996 are evident (see text).

(3 ppbv or less) are measured throughout the extensive cold region (not shown).

While maps constructed from gridded data provide a visually effective means of displaying observations, they can be deceptive in that the gridding process necessarily smooths out the contributions from individual measurements. PSCs are confined to the cold regions, which in early winter can be fairly localized; given the spacing of the UARS orbit tracks, short-lived PSCs of limited geographical extent may be represented in only one or two MLS data points, or may go undetected altogether. In Plate 2a the retrieved HNO₃ mixing ratios for June 2 are shown again, this time only at the MLS measurement locations, and in Plate 2b the UKMO temperatures are shown interpolated to these locations. Relatively low HNO₃ mixing ratios are recorded for a few points

where the temperatures are in the range 191–193 K. Although this area does register lower mixing ratios in the map of gridded data (Plate 1), the depression in HNO₃ in the low-temperature region is smaller than suggested by the individual measurements. In this study we mostly use averages of MLS data from several measurement locations rather than either individual observations (because of the noise associated with each retrieved profile; see section 2.2) or gridded data (because of the misrepresentation of small-scale features inherent in the gridding process). As discussed further below, the relationship between reductions in gas-phase HNO₃ and variations in temperature is also explored by binning the data with respect to temperature.

The vertical extent of severely depleted gas-phase HNO₃ and its relationship to low temperatures over the UARS

month are illustrated in Plate 3, which shows MLS HNO₃ as well as UKMO temperature and PV over the potential temperature range 420–585 K. The pink lines are contours of PV (scaled to yield similar values throughout the θ domain) that circumscribe the boundary of the vortex; the cyan and purple

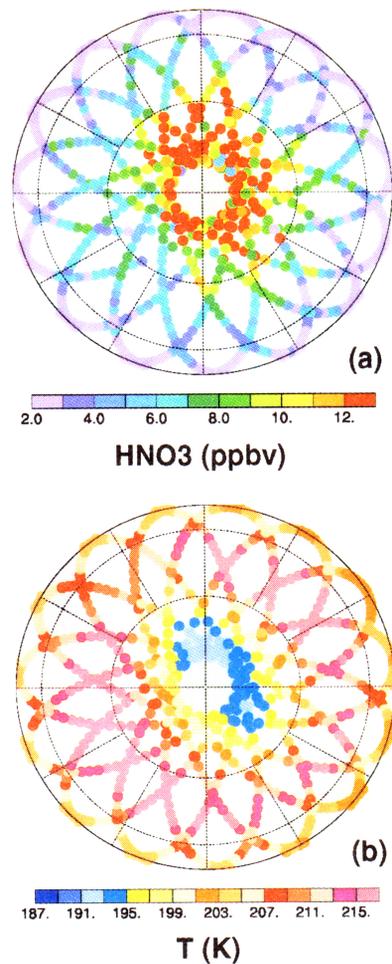
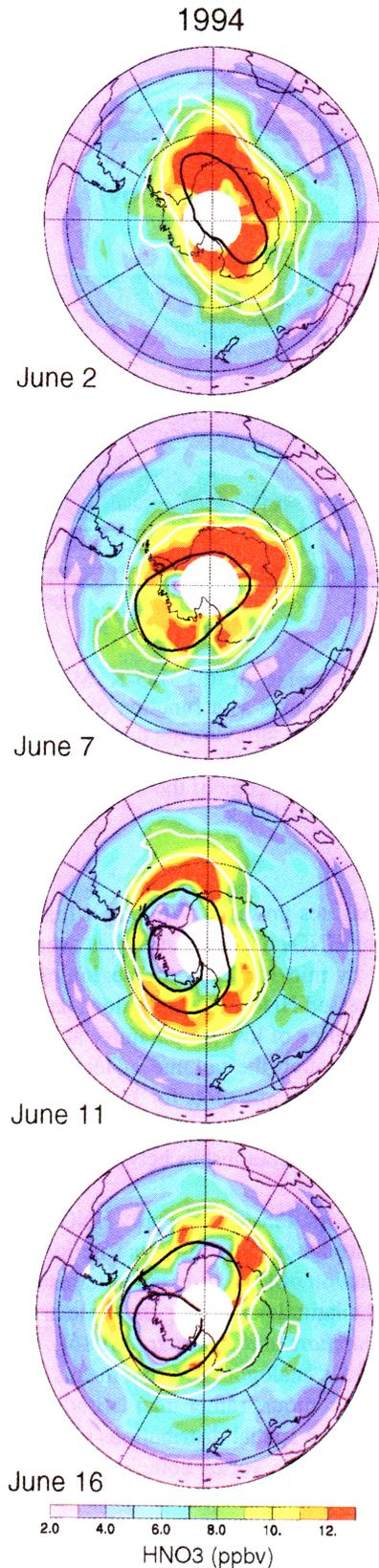


Plate 2. (a) The 465-K HNO₃ mixing ratios shown at the MLS measurement locations, with no horizontal binning or averaging performed, for June 2, 1994. (b) The 465-K UKMO temperatures for June 2, 1994, interpolated to the latitude/longitude coordinates of the MLS measurements.

lines demark the volumes within which the temperatures are below 192 K and the HNO₃ mixing ratios are below 4 ppbv, respectively. The surfaces of low HNO₃ along the perimeter of these plots are situated outside the vortex and are not associated with PSC formation. Temperatures in the upper

Plate 1. Maps of MLS HNO₃ (parts per billion by volume) for selected days during the 1994 early southern winter south viewing period, interpolated onto the 465-K potential temperature surface using UKMO temperatures. The maps are polar orthographic projections extending to the equator, with the Greenwich meridian at the top and with dashed black circles at 30°S and 60°S; blank spaces represent data gaps or bad data points. Superimposed in white are two contours of UKMO potential vorticity (PV): $-0.25 \times 10^{-4} \text{ K m}^2 \text{ kg}^{-1} \text{ s}^{-1}$ (to represent the approximate edge of the winter polar vortex at this level) and $-0.30 \times 10^{-4} \text{ K m}^2 \text{ kg}^{-1} \text{ s}^{-1}$ (a second contour to indicate the steepness of the PV gradient and thus the strength of the vortex). Superimposed in black are two contours of UKMO temperature: 195 K (the approximate existence threshold for type I polar stratospheric clouds (PSCs) and 188 K (the ice frost point).

portion of the vertical range are already well below 192 K by early June, but the bottom portion of the vortex does not cool below 192 K until June 8 (although 465-K temperatures are below 195 K at this time, as stated above). The HNO₃ depletion begins at the lowermost levels of the vortex and

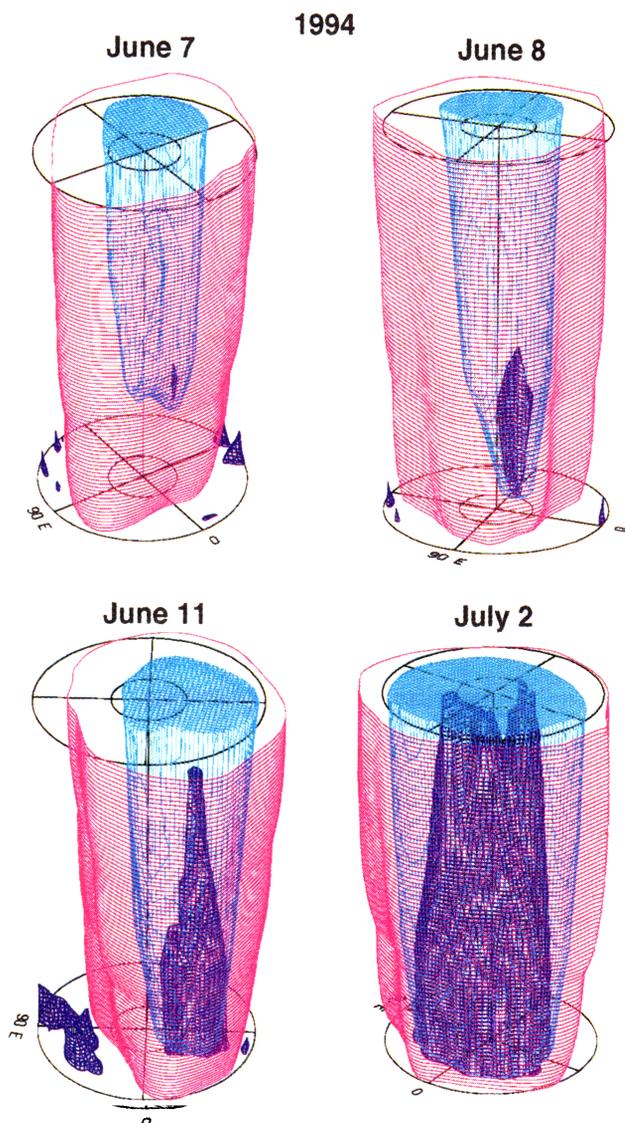


Plate 3. Plots illustrating the three-dimensional relationship between low temperature and severely reduced gas-phase HNO₃ inside the polar vortex for selected days during the 1994 early southern winter south viewing period. The vertical coordinate is potential temperature, ranging from 420 K to 585 K. Latitude circles are shown at 60°S and 80°S, and longitudes 0° and 90°E are labeled. The pink surface shows the 1.4 s⁻¹ contour of scaled PV (corresponding to the $-0.30 \times 10^{-4} \text{ K m}^2 \text{ kg}^{-1} \text{ s}^{-1}$ PV contour at 465 K [see Manney *et al.*, 1994]), which represents the approximate boundary of the polar vortex. The cyan surface shows the 192-K contour of UKMO temperature. The purple surface shows the 4-ppbv contour of MLS HNO₃. The surfaces of low HNO₃ along the perimeters of these plots are situated outside the vortex and are not associated with PSC formation; the geographic coverage of the HNO₃ plot is terminated at 60°S to minimize the presence of these surfaces. Different perspectives are used on each day to obtain the most unobstructed view of the vortex.

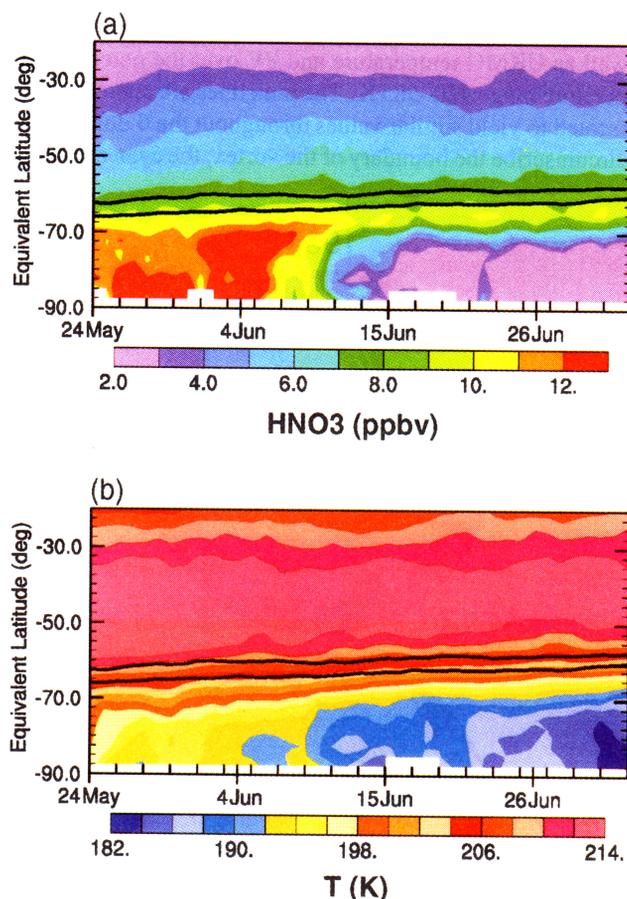


Plate 4. Time series of (a) 465-K MLS HNO₃ and (b) UKMO temperatures as a function of PV for the 1994 early southern winter south viewing period. PV is expressed in terms of equivalent latitude (see text). PV contours representing the approximate edge of the winter polar vortex (see Plate 1) are overlain in black.

gradually propagates upward. While still highly localized horizontally, the HNO₃ depletion stretches over a significant vertical extent by mid-June. By the beginning of July, the volume of depressed HNO₃ concentrations has expanded to nearly fill the volume of low temperatures at all but the top-most levels. Although the volume of temperatures below 188 K is smaller, the same plots for June 11 and July 2 using the 188-K temperature contour (not shown) indicate a substantially similar relationship between temperature and gas-phase HNO₃.

The evolution of 465-K gas-phase HNO₃ and UKMO temperatures throughout the observing period is shown in Plate 4 as a function of PV expressed in terms of equivalent latitude (i.e., the latitude enclosing the same area as a given PV contour [Butchart and Remsberg, 1986]). The quantities represented are analogous to zonal means but are calculated around PV contours rather than latitude circles. Both the temporal and the spatial distributions of reduced HNO₃ mixing ratios inside the vortex coincide closely with those of temperatures below 192 K. In contrast, the distribution of low HNO₃ does not correspond to that of temperatures below 195 K, the value commonly assumed as the threshold for PSC formation. The correlation in time between varia-

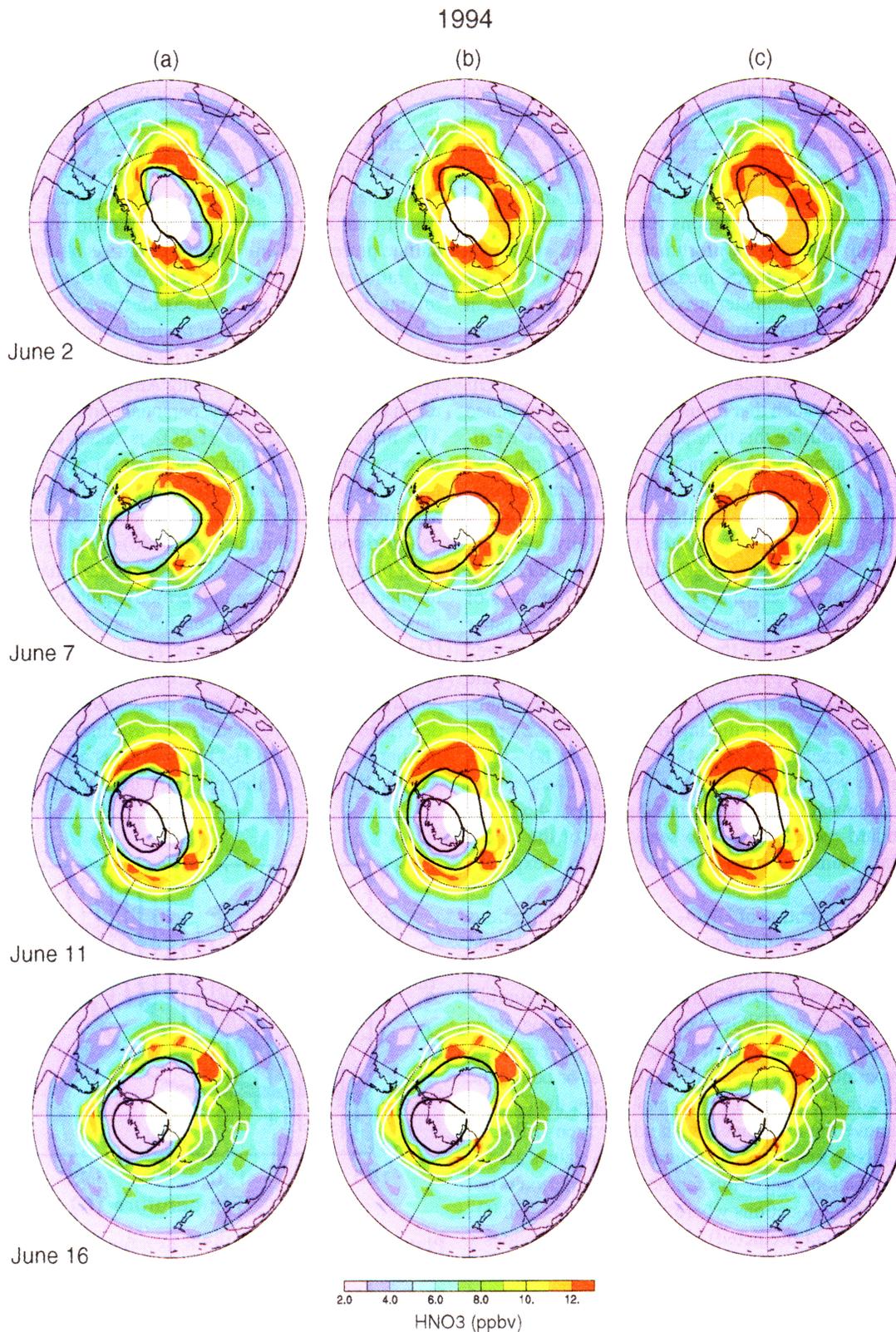


Plate 5. Maps at 465 K for the days during the 1994 early southern winter south viewing period shown in Plate 1, constructed from both MLS HNO₃ measurements and model calculations of the equilibrium vapor pressure of HNO₃ over (a) nitric acid trihydrate (NAT), using the formula of *Hanson and Mauersberger* [1988], (b) nitric acid dihydrate (NAD), using the formula of *Worsnop et al.* [1993], and (c) liquid ternary solutions, using the model of *Tabazadeh et al.* [1994b]. The maps are based on a H₂O mixing ratio of 4.5 ppmv and are constructed using the model values for HNO₃ inside the region where the temperature is between 186 K and 196 K and the MLS measurements of HNO₃ outside of this region.

tions in MLS HNO₃ mixing ratios and temperature fluctuations (noted for the 1995/1996 Arctic winter by *Santee et al.* [1996]) can also be seen in Figure 1. Again, there is a strong correspondence between the area of gas-phase HNO₃ loss and the area of temperatures in the range 190–192 K but only a weak correspondence between the area of gas-phase HNO₃ loss and the area of temperatures below 195 K.

The changes in the observed HNO₃ mixing ratios can be compared with those predicted by various models of PSC formation. These calculations require a value for the water vapor mixing ratio. Unfortunately, there are no simultaneous colocated UARS measurements of water vapor since the MLS radiometer used to measure it failed in April 1993 and the CLAES supply of cryogen was exhausted in May 1993. Although the UARS Halogen Occultation Experiment (HALOE) also measures water vapor, the HALOE observing pattern is such that no measurements are obtained poleward of 50°S during this time period. Therefore we estimate the water vapor abundance inside the vortex from 1992 MLS observations and from NASA ER 2 aircraft measurements made during ASHOE flights on June 1 and June 3, 1994, that penetrated into the vortex (K. Kelly, personal communication, 1996). MLS H₂O data for the 1992 early southern winter period are sparse, with only a few days of data in June and mid-July; they indicate initial water vapor mixing ratios of 4–5 ppmv throughout most of the vortex declining to less than 3 ppmv in the coldest regions near the pole by mid-July [*Santee et al.*, 1995]. The ER 2 measurements indicate water vapor mixing ratios of 5.0–5.2 ppmv along the flight track traversing 455–460 K and 60°–65°S (just inside the vortex). On the basis of these observations we assume a constant water vapor mixing ratio of 4.5 ppmv for the standard model runs. The sensitivity of the calculations to the assumed water vapor mixing ratio is explored below.

Maps of the distribution of HNO₃ for various assumptions about the controlling thermodynamics are constructed and compared with the data (Plate 5). We use the formula of *Hanson and Mauersberger* [1988] to calculate the equilibrium vapor pressure of HNO₃ over NAT at the geographic location of each MLS data point for which the interpolated UKMO temperature is less than 196 K (i.e., at or below the existence threshold for type I PSCs) and greater than 186 K (i.e., at or above the existence threshold for type II PSCs). To construct these maps, the measured HNO₃ values are retained at the locations where this temperature criterion is not satisfied (i.e., at the locations where sequestration of HNO₃ in type I PSCs is not predicted). This approach provides a picture of what MLS would have observed had NAT PSCs been forming according to the expression of *Hanson and Mauersberger* [1988]. Comparison of the maps in Plate 5a with those of MLS measurements in Plate 1 clearly shows that although temperatures are low enough to support NAT PSCs, they are not forming, at least not over spatial scales comparable to or larger than the $\sim 400 \times 100$ km horizontal resolution of the MLS measurements. Similarly, we use the formula of *Worsnop et al.* [1993] to calculate the equilibrium vapor pressure of HNO₃ over NAD and construct the

maps of Plate 5b. Although at first the spatial extent of low calculated mixing ratios is considerably larger than what is observed, by mid-June the NAD model provides fairly good agreement with the MLS HNO₃ data.

To determine the HNO₃ vapor pressure over ternary solutions, we use the aerosol physical chemistry model (APCM), described in detail by *Tabazadeh et al.* [1994b]. The APCM is an equilibrium model that calculates the partitioning of HNO₃ between the gas and the aerosol phase for given values of total HNO₃ and H₂O mixing ratio and total mass of sulfate present per unit volume of air. As above, a H₂O mixing ratio of 4.5 ppmv is used. We assume a total HNO₃ mixing ratio of 12 ppbv based on the MLS measurements taken inside the vortex prior to significant PSC activity. This value is consistent with in situ measurements of total reactive nitrogen (NO_y) obtained during ASHOE flights into the vortex in early June (D. Fahey, personal communication, 1996). The mass of condensed H₂SO₄ is taken to be 0.3 $\mu\text{g}/\text{m}^3$ to represent background (nonvolcanic) conditions. Again, this value is consistent with vortex measurements of sulfate aerosol volume from the early June ASHOE flights (J. Wilson, personal communication, 1996).

Application of the APCM assumes that the ternary droplets themselves are in thermodynamic equilibrium. Extensive comparisons between field observations and theory [*Drdla et al.*, 1994; *Carshaw et al.*, 1994; *Tabazadeh et al.*, 1994a; *Dye et al.*, 1996; *Del Negro et al.*, 1997] have confirmed that the growth of stratospheric aerosols is well described by ternary models that assume equilibrium partitioning of HNO₃ and H₂O between gas and aerosol phases. Although some of the particles may depart significantly

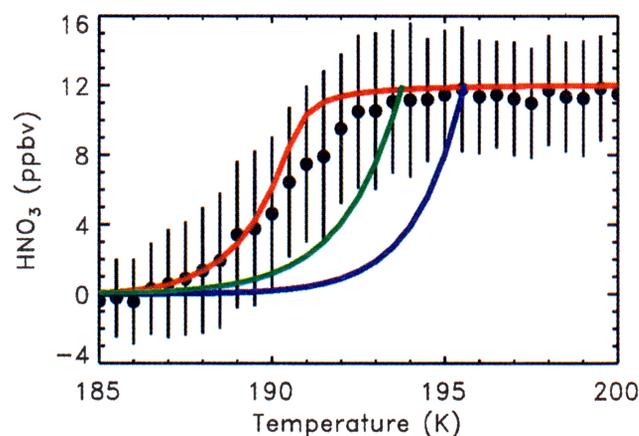


Plate 6. MLS HNO₃ at 465 K as a function of UKMO temperature for the 1994 early southern winter south viewing period. Approximately 8000 MLS data points inside the polar vortex (within the $-0.30 \times 10^{-4} \text{ K m}^2 \text{ kg}^{-1} \text{ s}^{-1}$ contour of UKMO PV) were binned into 0.5-K temperature intervals and averaged (black circles; standard deviations are indicated by error bars). Also shown are the equilibrium vapor pressures of HNO₃ over NAT (blue line), NAD (green line), and liquid ternary solutions (red line) calculated assuming a H₂O mixing ratio of 4.5 ppmv, a total HNO₃ mixing ratio of 12 ppbv, and a sulfate mass loading of 0.3 $\mu\text{g}/\text{m}^3$ (background conditions).

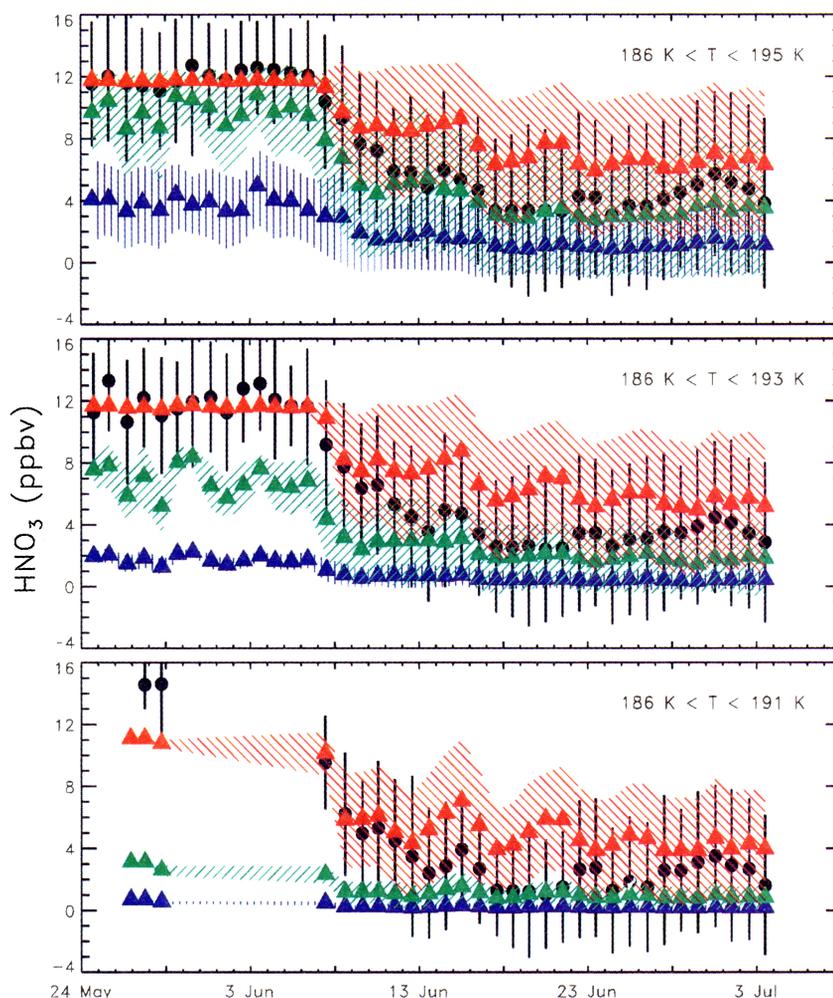


Plate 7. Time series (for the 1994 early southern winter south viewing period) of 465-K MLS HNO₃ (black circles) and calculated (see text for details) equilibrium vapor pressure of HNO₃ over NAT (blue triangles), NAD (green triangles), and ternary solutions (red triangles), averaged over the points enclosed within the annuli between the UKMO temperature contour of 186 K and the contours of 195, 193, and 191 K. The calculations were made using a total HNO₃ mixing ratio of 12 ppbv, a H₂O mixing ratio of 4.5 ppmv, and a sulfate mass loading of 0.3 $\mu\text{g}/\text{m}^3$ (background conditions). The standard deviations in the data averages are represented by the error bars, and the standard deviations in the model averages are indicated by the hatched areas.

from equilibrium, the overall aerosol volume matches that of the equilibrium ternary model very well, implying that the ternary system is not strongly influenced by the mesoscale temperature fluctuations [Murphy and Gary, 1995] that occur continuously in the stratosphere. Even under the rapid cooling conditions associated with lee wave events, while the smallest and the largest droplets are affected, the majority of the particles remain very close to the equilibrium composition [Meilinger et al., 1995]. Therefore in this study we have assumed that if only liquid particles are present in an air mass, then the HNO₃ uptake can be accurately represented by an equilibrium ternary model.

Maps constructed from the MLS data and the ternary solution calculations are shown in Plate 5c. Comparison of Plates 5a–5c shows that NAT particles are the most efficient, and ternary droplets the least efficient, in depleting HNO₃

from the gas phase. The ternary solution calculations are in much closer (albeit not perfect) agreement with the MLS data during the first days of PSC formation. Eventually, however, the ternary model underestimates the spread of HNO₃ loss. For intermediate days such as June 11, the best match to the data would probably be provided by a nonequilibrium mixture of NAD (Plate 5b) and ternary (Plate 5c) aerosols.

The relationship between observed HNO₃ and temperature for the entire UARS month is compared to predictions from the NAT, NAD, and liquid ternary solution models in Plate 6. For this plot all of the 465-K MLS measurements obtained inside the polar vortex (approximately 8000 data points) were binned into 0.5-K temperature intervals and averaged (black circles, with the standard deviations represented by error bars). Despite the large amount of scatter in

the MLS data, it is clear that neither the NAT (blue line) nor the NAD (green line) models match the average behavior of the MLS HNO₃ over most of the temperature range. In contrast, the ternary solution model (red line) exhibits a temperature dependence similar to that of the MLS data, although it predicts too little gas-phase HNO₃ loss for temperatures between 190 K and 193 K. All three models converge to match the observations obtained at the lowest temperatures, which were mainly concentrated toward the end of the month.

A comparison of the changes in the observed and predicted HNO₃ values over time is presented in Plate 7. Averages of 465-K MLS HNO₃ (black circles, with the standard deviations represented by error bars) are calculated from points enclosed within the annuli between the 186-K temperature contour and the 195-, 193-, and 191-K temperature contours. The 186-K lower limit is imposed because no allowance is made for the formation of water ice particles. Toward the end of the observing period, the 186-K lower bound eliminates most of the low-HNO₃ points in the vortex interior and leads to increases in the averaged mixing ratios. For each data point within these temperature contours the equilibrium vapor pressure of HNO₃ over NAT, NAD, and ternary solutions is calculated as described above in connection with Plate 5. Averages of the model results are also calculated within the different annuli and are depicted by blue (NAT), green (NAD), and red (ternary solution) triangles. The standard deviations in the model averages are indicated by the hatched areas.

The assumption of a NAT composition for PSCs leads to very low HNO₃ vapor concentrations that are inconsistent with the MLS observations from late May and early June. The NAD model predicts HNO₃ values that match the data closely during the latter half of the observing period; although the initial trend for NAD is similar to that seen in the data for the 195-K contour average, agreement with the data is much poorer for the averages at lower temperatures. In contrast, the behavior predicted for a ternary solution composition is in good agreement with the data during the early stages of HNO₃ depletion for all three temperature contour averages, with the best agreement attained at the lowest temperatures. Comparison of Plate 7 and Figure 2 shows that significant incorporation of HNO₃ into the ternary solution aerosols does not occur until the average temperatures drop below ~192 K.

Interpretation of the MLS data on the 465-K isentropic surface is potentially complicated by the processes of denitrification, whereby large PSC particles remove HNO₃ through gravitational settling out of the layer, and renitrification, whereby particles falling from above enhance HNO₃ through evaporation in the layer. Although various scenarios have been proposed in which denitrification takes place independent of the formation of water ice clouds [Salawitch *et al.*, 1989; Toon *et al.*, 1990b; Murphy and Gary, 1995], in general, the most intense denitrification is expected to be accompanied by dehydration. Since we have omitted from our analyses those regions where type II PSC formation is possible, the effects of severe denitrification have been largely eliminated. However, transport processes in the vortex lead

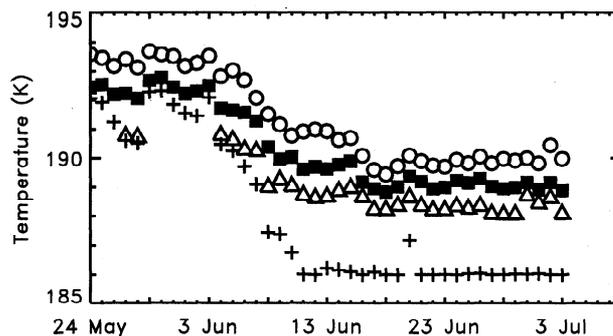


Figure 2. Time series (for the 1994 early southern winter south viewing period) of 465-K UKMO temperatures averaged over the annuli between the temperature contours of 186 K and 195 K (open circles), 193 K (solid squares), and 191 K (open triangles). Also shown are the minimum 465-K UKMO temperatures within those regions (plus signs; note that the minimum temperatures within the vortex as a whole, which are not shown, are much lower).

to dilution of the high HNO₃ concentrations in unperturbed regions as denitrified air is mixed in. To estimate the impact of denitrification on the total HNO₃ available in the PSC-free portions of the vortex, daily averages of HNO₃ are calculated from points located inside the vortex but outside the regions where HNO₃ mixing ratios are below 10 ppbv. These averaged HNO₃ values (not shown) exhibit a decrease of ~1.5 ppbv over the course of the observing period. Slightly larger decreases are obtained when averages are calculated over regions bounded by smaller HNO₃ mixing ratios (e.g., ~2.5 ppbv for a HNO₃ cutoff value of 4 ppbv). Because of the long timescale for HNO₃ photolysis during polar winter in the lower stratosphere [e.g., Austin *et al.*, 1986], the decline in these averaged HNO₃ values probably arises mainly from denitrification. The ternary model is very sensitive to the total amount of HNO₃. Lower abundances of available HNO₃ (e.g., ~6–8 ppbv, rather than 12 ppbv as in Plate 7; not shown) brought about by denitrification would improve the agreement between the ternary model and the observations toward the end of the observing period. However, the ternary droplets themselves are relatively small in size [e.g., Toon *et al.*, 1990a] and are unlikely to cause significant denitrification, especially not in the short period of time over which the best agreement with the measurements switches from the ternary to the NAD model (see Plate 7). Unlike the ternary model, the NAD predictions are not strongly dependent on the total HNO₃. Although the ternary model could be made to match the data well throughout the entire observing period by invoking either rapid denitrification or rapid HNO₃ photolysis, a more likely explanation of the behavior depicted in Plate 7 involves a change of phase from ternary solutions to NAD.

Temperatures in the layers above 465 K are low enough for type I and, in the latter half of the study period, type II PSC formation. Evidence of the vertical redistribution of HNO₃ through evaporation of falling PSC particles has been seen in both hemispheres [Arnold *et al.*, 1989; Hübler *et*

al., 1990; Hofmann and Deshler, 1991; Kawa *et al.*, 1992; Tuck *et al.*, 1995]. In these cases, however, HNO₃ enhancement was observed at much lower levels in the atmosphere. Renitrification requires evaporation to halt particle sedimentation and liberate gas-phase HNO₃. Temperatures at 465 K remain below the NAT existence threshold in much of the vortex throughout this UARS month. Therefore any PSC particles that form at higher altitudes and are large enough to undergo gravitational settling would likely pass through the 465-K level without evaporating.

Another process affecting HNO₃ abundances in the lower stratosphere is diabatic descent. Results from a three-dimensional Lagrangian transport calculation (not shown) similar to those described by Manney *et al.* [1995a] indicate that descent leads to an increase in vortex-averaged HNO₃ at 465 K of ~0.5 ppbv over the 30 days between June 2 and the end of the observing period. However, the calculations show essentially no change in HNO₃ when averages are limited to the low-temperature areas. The differences between the two sets of averages reflect the fact that descent during early southern winter is stronger along the vortex edge than in its interior [Manney *et al.*, 1994], where the cold regions are predominantly confined. In any case, vertical transport produces only small increases in the HNO₃ mixing ratios at 465 K during this UARS month. Since these increases are offset by the effects of denitrification and HNO₃ photolysis discussed above, we do not consider them further.

Of greater importance in interpreting the MLS HNO₃ measurements are the uncertainties associated with various model parameters. One means of assessing these uncertainties is to explore a range of parameter values in performing the model calculations of Plate 7. Plate 8, discussed below, illustrates the sensitivity of the model calculations to the assumed water vapor mixing ratio, the temperature, and the assumed sulfate aerosol loading. For all cases shown, the averages are calculated from points enclosed within the annulus between the 186-K and 195-K temperature contours using a total HNO₃ mixing ratio of 12 ppbv; these results can be compared with those of the top panel of Plate 7. In general, the model calculations exhibit much less sensitivity for averages computed within the 191-K temperature contour.

Model calculations were performed for a series of water vapor mixing ratios using the nominal UKMO temperatures and background aerosol conditions. Results for the extremes of 3 and 6 ppmv are shown in Plate 8a since they bracket the behavior of the predicted HNO₃. Although the MLS data indicate water vapor mixing ratios on the order of 3 ppmv by mid-July, they do not support such extensive water vapor loss in the late May/early June time frame [Santee *et al.*, 1995]. Nevertheless, the top panel of Plate 8a shows that the agreement of both the NAT and NAD models with the data could be significantly better than that of the standard run if we have overestimated the water vapor abundance. In the case of water mixing ratios larger than those of our standard run, the bottom panel of Plate 8a shows even better agreement between the data and the ternary solution model.

We also investigate the influence of possible tempera-

ture biases on the model results. Differences between the UKMO analyses and radiosonde observations at 50 hPa at high southern latitudes during May and June are typically within about ± 1 K, with even smaller systematic biases (< 0.5 K) indicated [Manney *et al.*, 1996]. Varying the temperature by ± 1 K from the nominal UKMO value at each point yields results (not shown) in which the ternary and NAD models exhibit a similar level of agreement with the data, but the NAT model still significantly overestimates the HNO₃ depletion throughout the observing period. Results from the more extreme case of varying the temperature by ± 3 K from the nominal UKMO value at each point are shown in Plate 8b (for a total water vapor mixing ratio of 4.5 ppmv and background aerosol conditions). Note that shifting the temperature by this amount at each point causes a different set of measurements to be enclosed within the annulus between the 186-K and 195-K temperature contours, resulting in data averages much different from those of Plate 7. Although differences between UKMO and radiosonde temperatures of this magnitude occur only occasionally during this time period, Plate 8b shows that if the UKMO temperatures are systematically biased high by 3 K (top panel), then the discrepancy between the data and the NAT and NAD models is greater than indicated by the standard run, whereas the ternary model provides a fairly good match to the data throughout the entire month; if the UKMO temperatures are systematically biased low by 3 K (bottom panel), then the agreement of NAT with the data is improved over that of the standard run, particularly during the latter half of the period. Note that comparisons of UKMO and radiosonde observations [Manney *et al.*, 1996] do not support the existence of systematic biases larger than ~ 1 K at any time during the southern hemisphere winter and suggest that systematic biases are even smaller in May and June.

In Plate 8c the sulfate mass is varied between 3 and 30 $\mu\text{g}/\text{m}^3$ (using the nominal UKMO temperatures and a total water vapor mixing ratio of 4.5 ppmv). Although unaffected by the sulfate aerosol loading, the NAT and NAD model curves are retained for consistency. Enhancement of the total sulfate mass leads to greater uptake of gas-phase HNO₃, as shown previously by Carslaw *et al.* [1994] and Tabazadeh *et al.* [1994b]. The bottom panel of Plate 8c represents strongly perturbed volcanic conditions; the top panel shows that if we have underestimated the residual effects from the Mount Pinatubo eruption, then the agreement of the ternary model with the data could be worse at the beginning of the observing period than indicated by the standard run.

Finally, in Plate 8d we explore the potential impact on the model comparisons of bias errors in the MLS HNO₃ measurements. The model calculations are made using the standard values (H₂O = 4.5 ppmv, total HNO₃ = 12 ppbv, sulfate mass = 0.3 $\mu\text{g}/\text{m}^3$) as in Plate 7, but the MLS data are shifted by ± 2 ppbv at each measurement location. Plate 8d shows that if the HNO₃ retrievals interpolated to 465 K are biased high by 2 ppbv (top panel), then the agreement with the NAD model is considerably better than indicated by the results of Plate 7. On the other hand, a low bias of 2 ppbv in

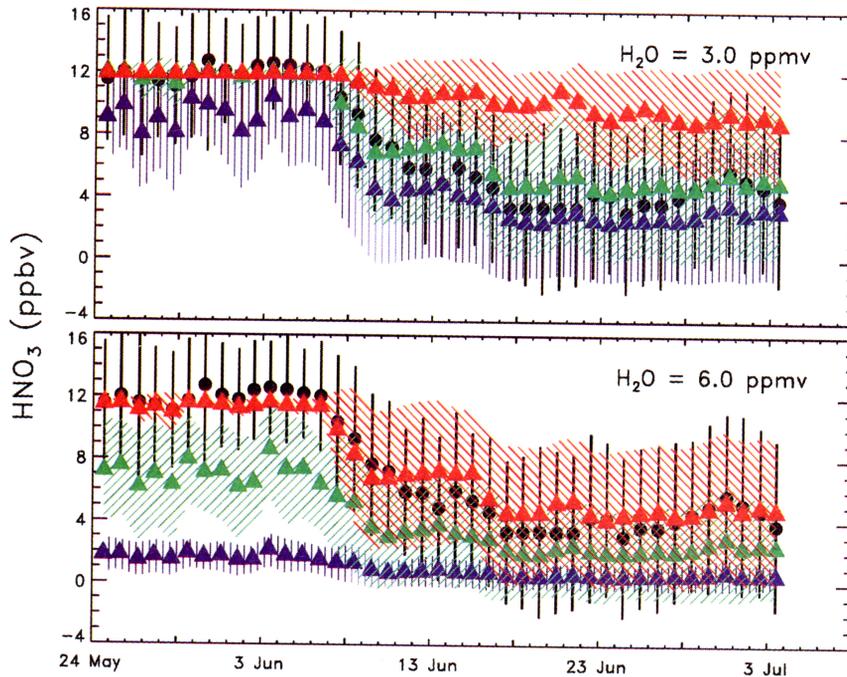


Plate 8a. Time series, as in Plate 7, to illustrate the sensitivity of the HNO₃ vapor pressure calculations to the assumed water vapor mixing ratio, which is varied between 3 and 6 ppmv. The averages are calculated from points enclosed within the annulus between the 186-K and 195-K temperature contours using a total HNO₃ mixing ratio of 12 ppbv.

the MLS data (bottom panel) would point to an even greater discrepancy between the data and the NAT and NAD models. Further validation efforts to quantify any bias errors in the MLS HNO₃ data are currently underway.

That the data are matched better by a ternary solution model during the initial stages of PSC formation but as time

progresses are matched better by a NAD model is consistent with a scenario proposed by *Tabazadeh et al.* [1994a, 1995, 1996] based on analysis of aircraft data. According to this scenario, ternary solutions form in air masses in which the background sulfate aerosols are liquid (e.g., in the first cooling cycles) and the temperatures have recently dropped a few

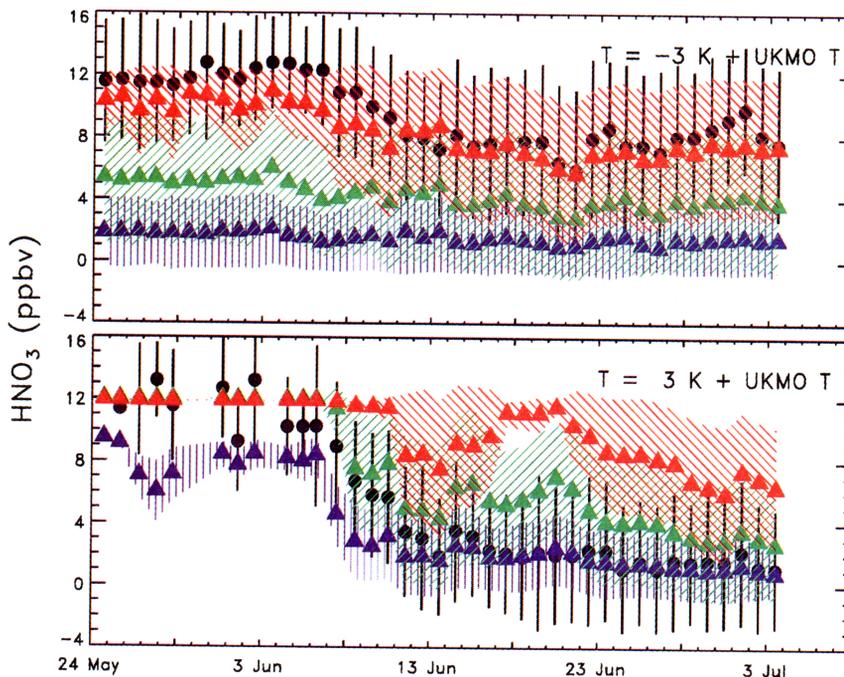


Plate 8b. Same as in Plate 8a, but for the temperature, which is varied by ± 3 K from the nominal UKMO temperature at each measurement location.

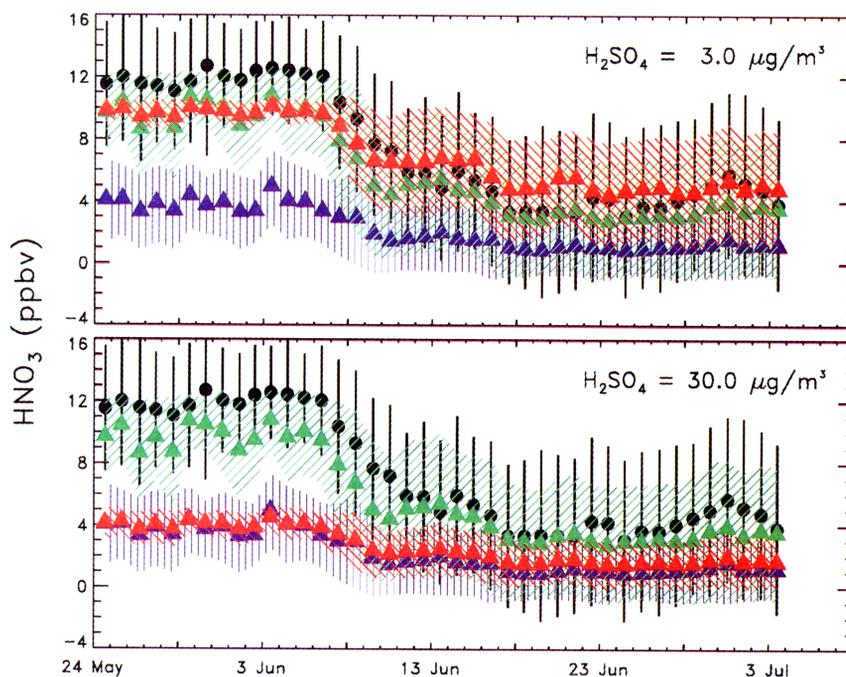


Plate 8c. Same as in Plate 8a, but for the assumed total sulfate mass, which is varied between 3 and 30 $\mu\text{g}/\text{m}^3$.

degrees below the NAT condensation point. After exposure to low temperatures for more than ~ 1 day, NAT or NAD particles begin to nucleate, and a mixed cloud results while a transfer of HNO₃ vapor from the ternary solution droplets to the more stable crystalline particles takes place, with eventual complete conversion. That is, a nonequilibrium situation arises after some of the ternary droplets freeze, and

the HNO₃ vapor pressure, initially at the ternary limit, gradually decreases as the HNO₃ is transferred to the thermodynamically more favorable NAT or NAD particles. The MLS observations (see Plate 7) suggest that the first crystalline structure to form is NAD, the preferential nucleation and persistence of which was first propounded for the polar stratosphere by Worsnop *et al.* [1993]. In addition, the

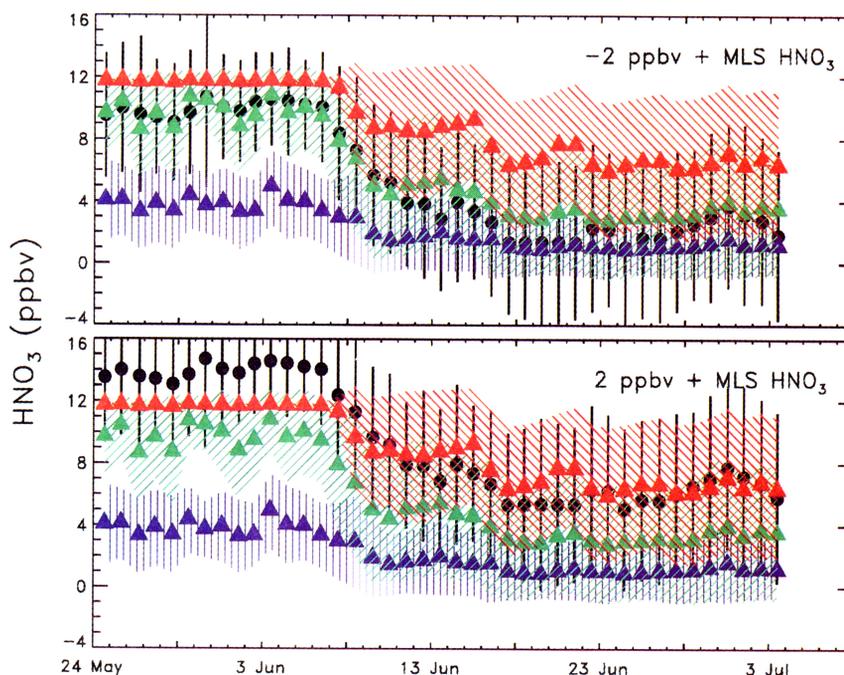


Plate 8d. Time series, as in Plate 7, to explore the potential impact on the model comparisons of bias errors in the MLS HNO₃, which is varied by ± 2 ppbv from the nominal mixing ratio at each measurement location. Model calculations are made using the standard parameter values as in Plate 7.

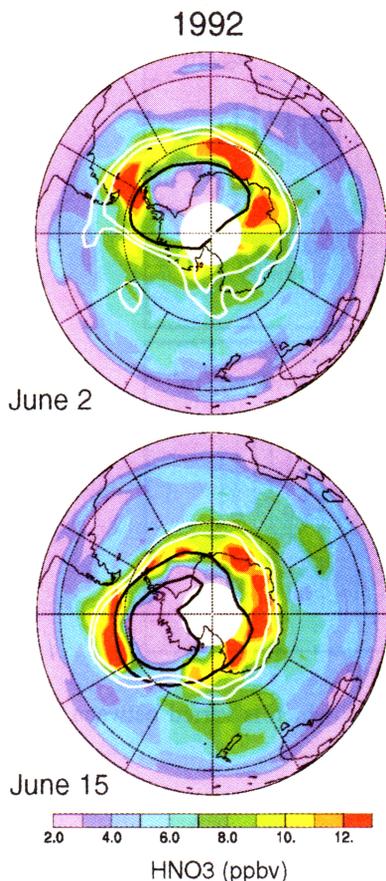


Plate 9. Same as in Plate 1, but for the 1992 early southern winter south viewing period.

MLS data indicate that the conversion of ternary solutions to NAD occurs more slowly (spanning at least several days) than suggested by *Tabazadeh et al.* [1996].

This scenario is also supported by the analyses of *Larsen et al.* [1996, 1997]. They examined balloon-borne backscatter measurements from several Arctic winters, along with

back trajectory calculations, and found that ternary solutions were observed predominantly in air masses that were in the process of undergoing relatively fast cooling, whereas most solid type Ia PSCs exhibiting the characteristics of NAT had aged at temperatures below the NAT condensation point for more than 1–2 days. In addition, *Dye et al.* [1996] reported that particle volumes observed at the lowest temperatures during the ASHOE PSC encounter were not entirely consistent with an equilibrium ternary model, and they suggested a mixture of ternary droplets with NAT or other solid phases. In section 5 we use trajectory calculations to further investigate the evolution of the PSC composition.

4.2. The 1992 Early Southern Winter

At the start of the 1992 early southern winter south viewing period on June 2 (Plate 9), HNO₃ abundances are high in the portion of the vortex where temperatures are above ~ 195 K, due to diabatic descent as discussed above. However, a deficit in gas-phase HNO₃ is apparent inside most (but not all) of the region where temperatures are below ~ 195 K, more or less conforming to the pattern expected for crystalline type I PSCs. These observations are in contrast to the large HNO₃ mixing ratios observed throughout most of the vortex, including the low-temperature region, in early June 1994 (see Plate 1 and Figure 1). Operational constraints precluded the acquisition of additional data during 1992 until June 15 (Plate 9), by which time temperatures over a broad region have dropped below 188 K, and HNO₃ maps for the two years appear substantially similar.

The difference in the behavior of MLS HNO₃ during early June in 1992 as compared with 1994 is probably related to the sulfate aerosol loading following the eruption of Mount Pinatubo [e.g., *Tabazadeh et al.*, 1994b; *Carslaw et al.*, 1994]. *Russell et al.* [1996] have shown that the stratospheric particle optical depth in April 1992 at 60°S–70°S was enhanced by roughly a factor of 20–30 over pre-Pinatubo levels. In Plate 10 we show HNO₃ maps con-

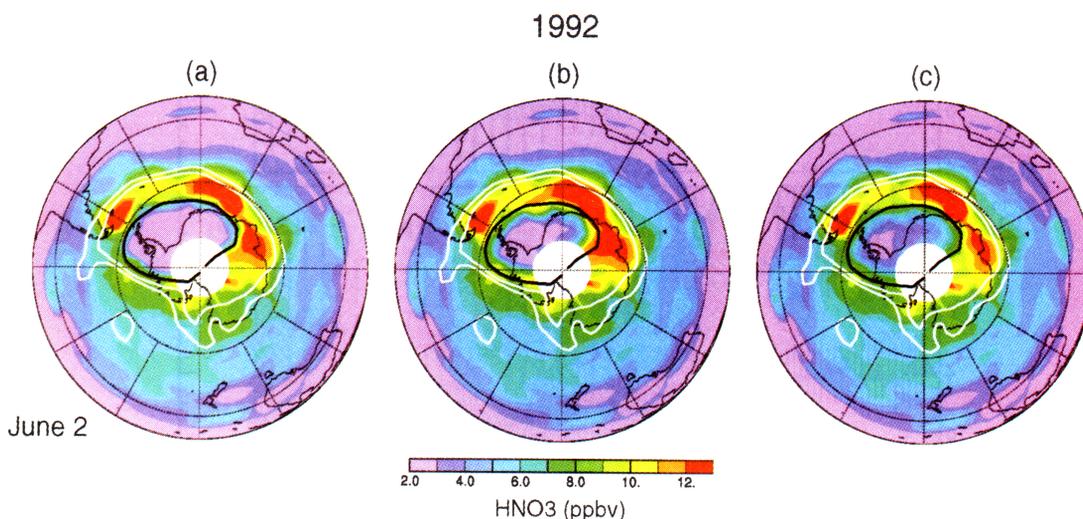


Plate 10. Same as in Plate 5 for (a) NAT, (b) NAD, and (c) liquid ternary solutions (using a total sulfate mass of $10 \mu\text{g}/\text{m}^3$), for June 2, 1992.

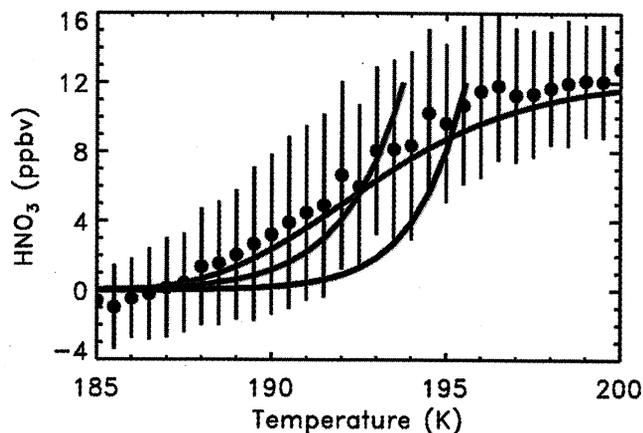


Plate 11. Same as in Plate 6, but for the 1992 early southern winter south viewing period. The averages represent approximately 6300 MLS data points inside the polar vortex. The ternary model curve was calculated using a sulfate mass loading of $10 \mu\text{g}/\text{m}^3$ to account for the Pinatubo enhancement.

constructed for June 2 from the MLS data and the various PSC composition models using a H₂O mixing ratio of 4.5 ppmv and a total HNO₃ mixing ratio of 12 ppbv, as discussed in connection with Plate 5, and a total sulfate mass of $10 \mu\text{g}/\text{m}^3$ to represent the Pinatubo enhancement. Although the ternary model significantly underestimates the HNO₃ uptake for sulfate mass loadings of $0.3 \mu\text{g}/\text{m}^3$ and $3 \mu\text{g}/\text{m}^3$ (not shown), it agrees fairly well with the data for a sulfate mass of $10 \mu\text{g}/\text{m}^3$. However, the models constrained by the thermodynamics of NAD and NAT also agree reasonably well with the observed HNO₃. The single MLS observation day at the beginning of the 1992 early southern winter study period does not allow us to definitively distinguish between any of the compositions. However, the temperature dependence of the HNO₃ measurements obtained inside the polar vortex during the course of the observing period, shown in Plate 11, indicates that a ternary solution model using a sulfate mass of $10 \mu\text{g}/\text{m}^3$ predicts reductions in gas-phase HNO₃ similar to those observed by MLS.

4.3. The 1993 Early Southern Winter

As in 1994, MLS obtained a complete set of daily limb scans during the 1993 early southern winter south looking period (May 29 to July 7; selected days shown in Plate 12). Unlike in 1994, however, in 1993, MLS did not resume south viewing in time to capture the incipient stage of PSC formation. As indicated in Figure 1, the areal extent of HNO₃ mixing ratios below 4 ppbv measured on the first day of the 1993 observing period is not achieved in 1994 until approximately June 9, midway through that observing period. The relationship between low HNO₃ and low temperature for these days is different for the two years: In 1993, the areal extent of low HNO₃ values is comparable to that of temperatures below 192 K, but in 1994 it is much smaller. In fact, in 1994 the areal extent of low HNO₃ remains smaller than that of temperatures below 190 K for several days. A possible explanation for this difference is the lingering aftereffects in 1993 of

the eruption of Mount Pinatubo. *Russell et al.* [1996] have shown that the stratospheric particle optical depth in April 1993 at 60°S to 70°S was still enhanced by roughly a factor of 10 over pre-Pinatubo levels. Equilibrium calculations

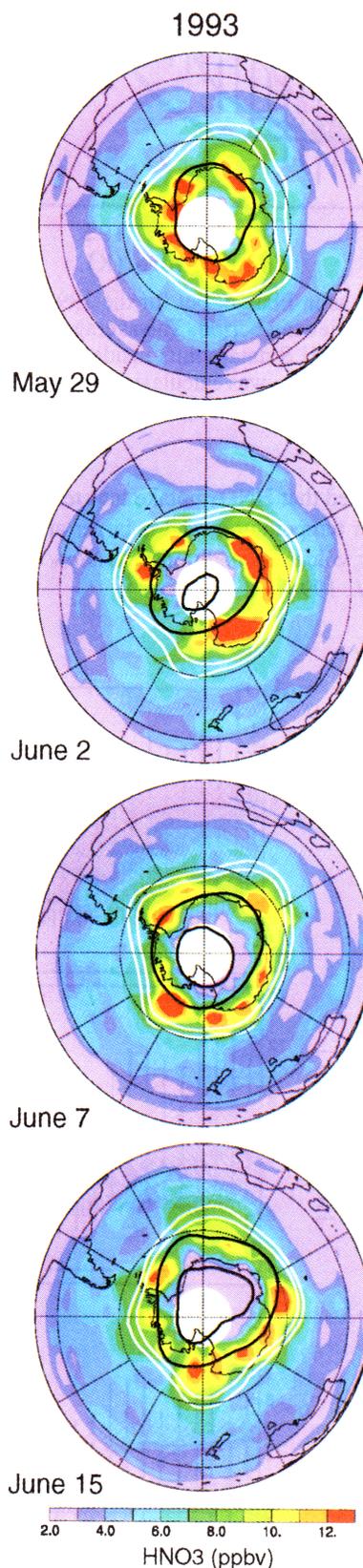


Plate 12. Same as in Plate 1, but for the 1993 early southern winter south viewing period.

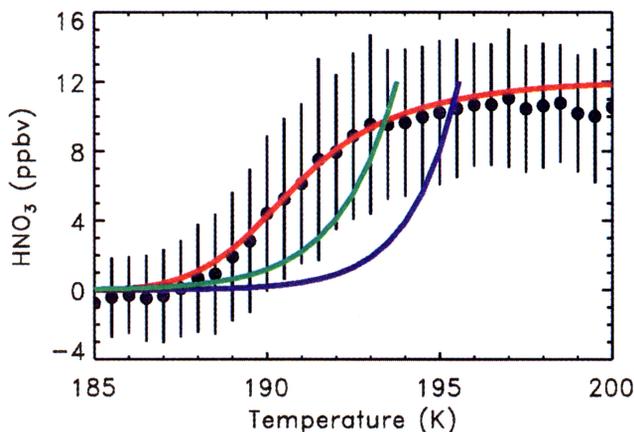


Plate 13. Same as in Plate 6, but for the 1993 early southern winter south viewing period. The averages represent approximately 8500 MLS data points inside the polar vortex. The ternary model curve was calculated using a sulfate mass loading of $3 \mu\text{g}/\text{m}^3$ to account for the residual Pinatubo enhancement.

have demonstrated that HNO₃ uptake into ternary solutions occurs at higher temperatures, and is larger at a given temperature, under volcanic conditions [Carlsaw *et al.*, 1994; Tabazadeh *et al.*, 1994b].

The HNO₃/temperature relationship inside the vortex for this observing period (Plate 13) shows that a ternary solution model using a total sulfate mass of $3 \mu\text{g}/\text{m}^3$ to account for the residual Pinatubo enhancement agrees well with the MLS data. In Plate 14 the changes over time in the observed HNO₃ mixing ratios are compared with those predicted by the various PSC composition models (using a H₂O mixing ratio of 4.5 ppmv, a total HNO₃ mixing ratio of 12 ppbv, and a total sulfate mass of $3 \mu\text{g}/\text{m}^3$). As was seen for the 1994 data, the NAT model severely underestimates the observed HNO₃ abundances, while the NAD model calculations match the data closely in the latter half of the observing period and the ternary model provides excellent agreement initially.

Comparison of Plate 14 with Plate 7 reveals that the ternary calculations provide a good description of the data after PSC formation begins for a much longer duration in 1993 than in 1994. This is probably related to the significant differences in composition between ternary solutions formed under volcanic conditions and those formed under background aerosol conditions [Tabazadeh *et al.*, 1994b]. These compositional differences during volcanically perturbed periods may give rise to changes in the NAT or NAD saturation ratio that hamper the formation of crystalline PSCs from the

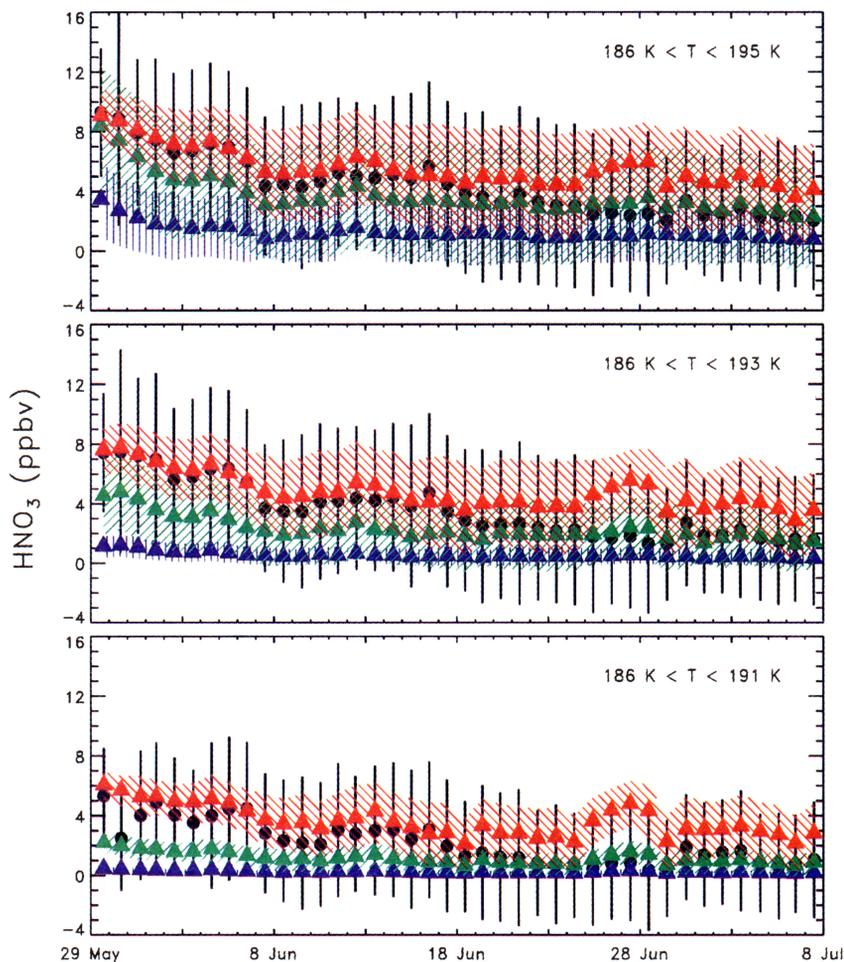


Plate 14. Same as in Plate 7, but for the 1993 early southern winter south viewing period, using a total HNO₃ mixing ratio of 12 ppbv, a H₂O mixing ratio of 4.5 ppmv, and a total sulfate mass of $3 \mu\text{g}/\text{m}^3$.

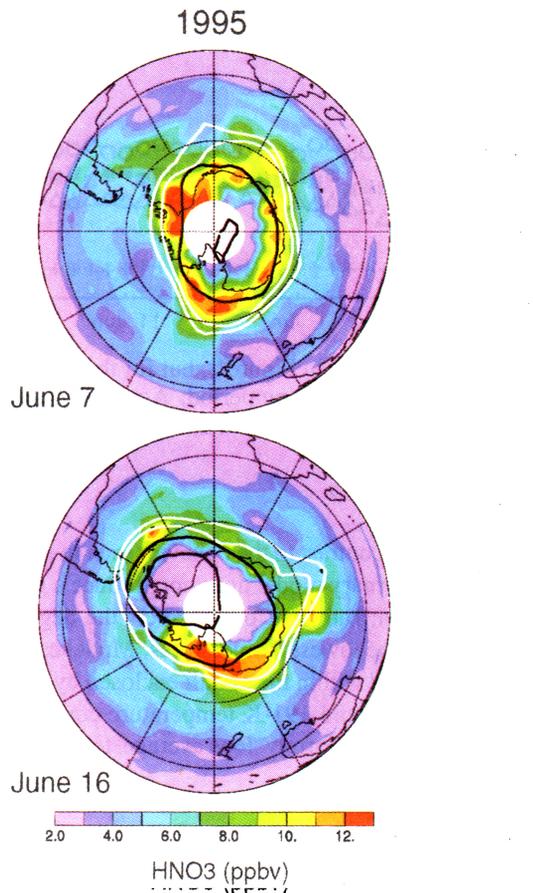


Plate 15. Same as in Plate 1, but for the 1995 early southern winter south viewing period.

ternary solutions. Plate 14 suggests, however, that the conversion to NAD did eventually take place in 1993.

4.4. The 1995 Early Southern Winter

The first and last days for which MLS data are available during the 1995 early southern winter are June 7 and June 16, respectively (Plate 15). Figure 1 shows that in 1995 temperatures have been below 190 K since the beginning of June, and by the first day of MLS observations there is already a substantial area of low HNO₃ mixing ratios. Again, MLS has not observed the onset of PSC formation. Although the NAD model produces HNO₃ depletion over too broad an area, especially for the first few days of measurements, it still provides the best agreement with the MLS data (not shown).

4.5. The 1996 Early Southern Winter

During the 1996 early southern winter the lower stratosphere was substantially colder than average [*Climate Prediction Center*, 1996], with 465-K minimum temperatures continuously below 195 K after May 7 and 188 K after May 22. The earlier onset and greater area of very low temperatures in the 1996 early winter is evident in Figure 1. Because of UARS power limitations, MLS measurements are not available during this period until May 22, by which time PSC formation has already begun (Plate 16). Although the progression of HNO₃ depletion appears similar for 1996 and

1994, its relationship with low temperature is not: Initially, the extent of the reduction in gas-phase HNO₃ within the regions of lowest temperature is smaller in 1996 than in 1994. In fact, Figure 1 shows that despite the very low temper-

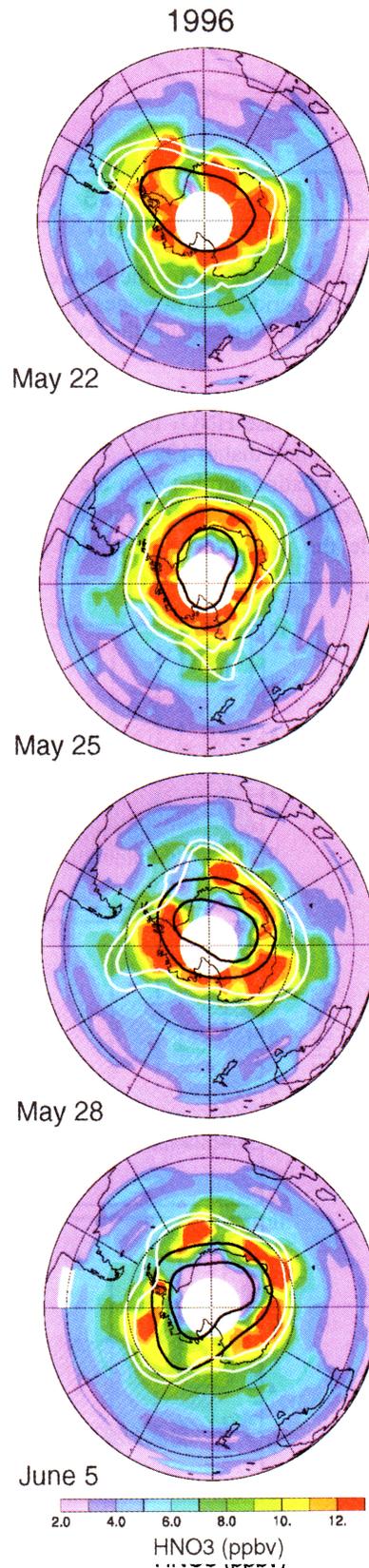


Plate 16. Same as in Plate 1, but for the 1996 early southern winter south viewing period.

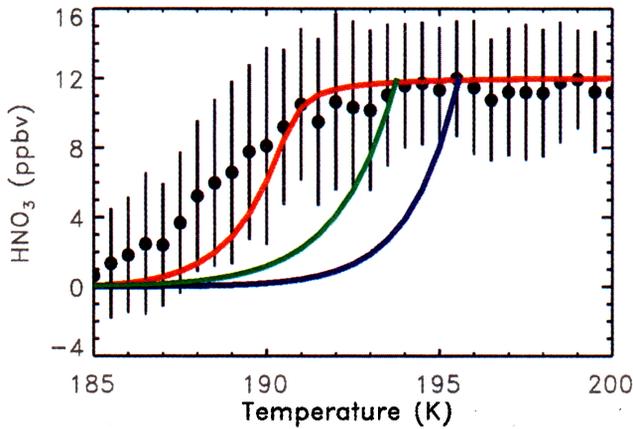


Plate 17. Same as in Plate 6, but for the 1996 early southern winter south viewing period. The averages represent approximately 2800 MLS data points inside the polar vortex.

atures in 1996, gas-phase HNO₃ depletion proceeds at an initial rate less than half that of 1994.

As before, the relationship between the observed HNO₃ and the temperature is examined for all measurements obtained inside the polar vortex over the UARS month

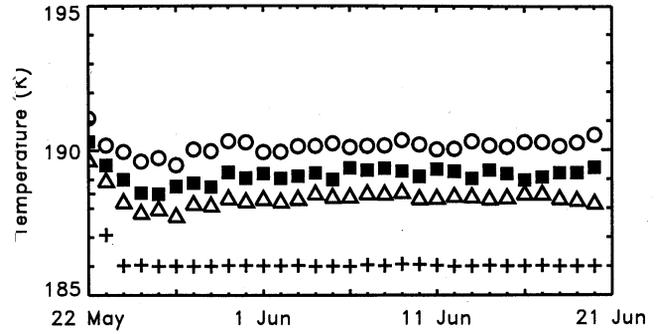


Figure 3. Same as in Figure 2, but for the 1996 early southern winter south viewing period. Again, the minimum temperatures within the vortex as a whole, which are not shown, are much lower.

(Plate 17). In contrast to the behavior seen in previous years, in 1996 the ternary model does not produce a good match to the MLS data; rather, it overpredicts the degree of HNO₃ depletion at the lowest temperatures. Model calculations are also performed within the annuli enclosed by the 186-K and the 195-, 193-, and 191-K temperature contours. Figure 3 illustrates the uniformly low average temperatures within

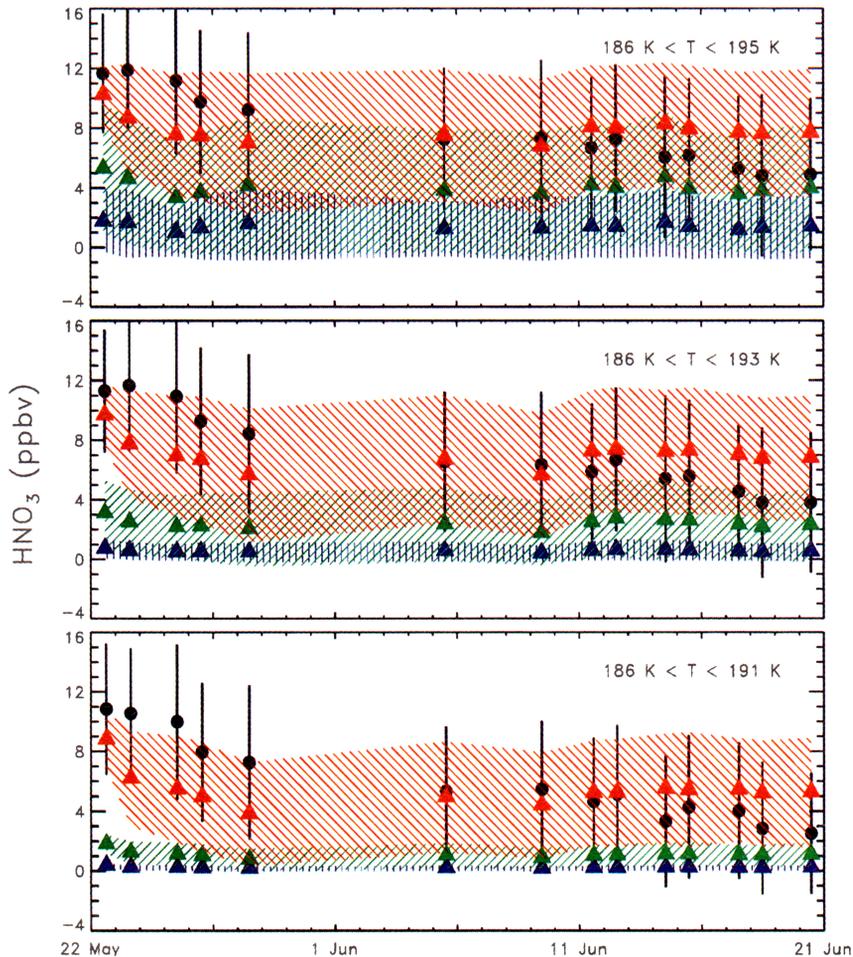


Plate 18a. Same as in Plate 7, but for the 1996 early southern winter south viewing period, using a HNO₃ mixing ratio of 12 ppbv, background aerosol conditions, and a H₂O mixing ratio of 4.5 ppmv.

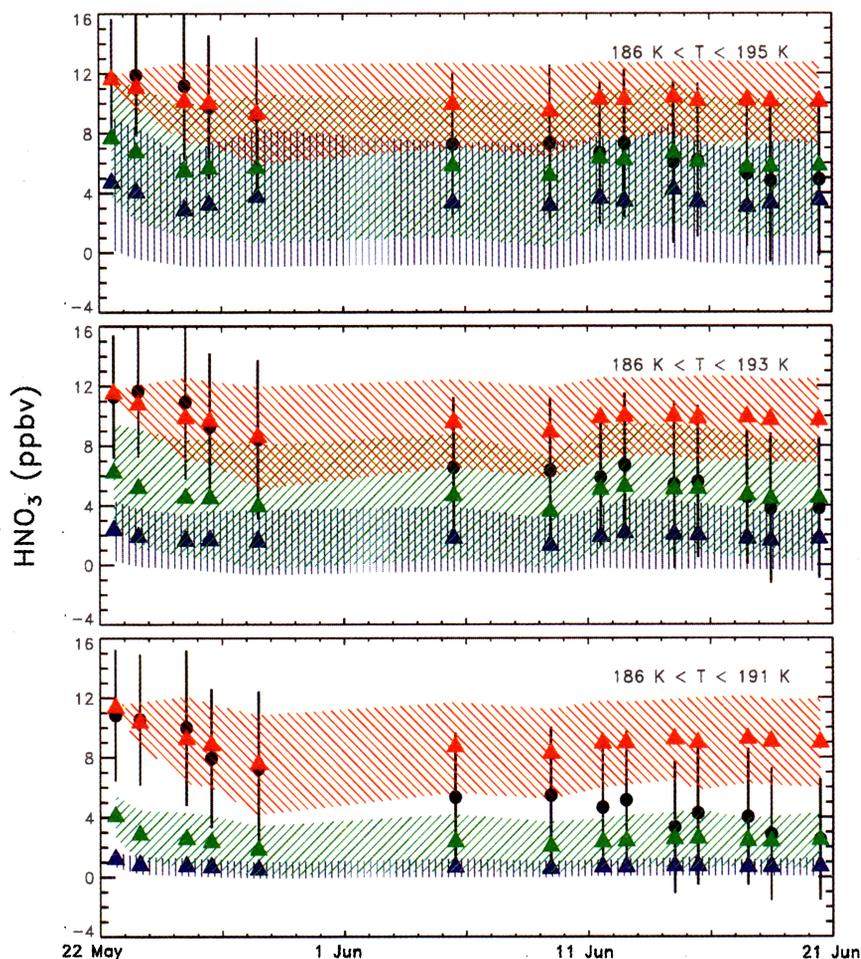


Plate 18b. Same as in Plate 18a, but for a H₂O mixing ratio of 3.0 ppmv.

these regions throughout the observing period. In Plate 18a we show the time series of observed HNO₃ mixing ratios averaged within the low-temperature regions along with those predicted by the various PSC composition models using the standard values of H₂O=4.5 ppmv, HNO₃=12 ppbv, and H₂SO₄=0.3 μg/m³. Again, in contrast to the results for earlier years, in 1996 the ternary solution model initially overestimates the degree of HNO₃ uptake. There is some indication in Plate 18a that the agreement with the data improves by the middle of the month for a ternary composition; this may be significant or it may simply be a coincidence since the MLS data exhibit a decreasing trend throughout this period.

It is unlikely that nonequilibrium conditions can explain the 1996 MLS observations. As discussed in section 4.1, even under conditions of rapid cooling liquid ternary solutions maintain a nearly equilibrium composition. A nonequilibrium mixture of NAT or NAD particles and ternary droplets is possible, but in this case the resulting HNO₃ vapor pressure would be between those of the solid and liquid phases. That is, if a majority of the particles are liquid, then gas-phase HNO₃ depletion would be at least as large as that due to HNO₃ uptake into ternary solutions; the presence of some solid particles (over which the HNO₃ vapor pres-

sure is lower) would lead to even greater gas-phase HNO₃ loss. Since the observed HNO₃ vapor pressure in Plate 18a is higher than that predicted by either NAT/NAD or ternary models, a cloud of mixed NAT/NAD and ternary phases can be ruled out.

To explore the possibility that the lack of agreement between the ternary model and the data arises because the assumed abundance of available HNO₃ (12 ppbv) is too low, we have run the ternary model with the total HNO₃ mixing ratio set to 15 ppbv (not shown). While a slightly better fit to the data is achieved for the 195-K contour average, improvements for the averages at lower temperatures are negligible. Results from a second sensitivity test in which the H₂O mixing ratio is reduced to 3 ppmv, while the HNO₃ mixing ratio is kept at 12 ppbv, are shown in Plate 18b. This case yields behavior more akin to that seen in previous years, with an excellent match to the data provided by the ternary model during the initial HNO₃ depletion and by the NAD model thereafter.

Although the HNO₃ and H₂O abundances utilized in Plate 18b allow a good match to the MLS data to be achieved, they may not be representative of actual stratospheric conditions. As discussed in section 4.1, water vapor mixing ratios at the beginning of the winter season would

typically be expected to be in the range of 4–5 ppmv. Values of 3 ppmv or less have been measured inside the vortex in August and September [Kelly *et al.*, 1989; Santee *et al.*, 1995], but these measurements followed the period of extensive dehydration due to the sedimentation of type II (water ice) PSCs. Sufficient time has not elapsed since the onset of low temperatures for significant dehydration to have occurred by May 22, but even if it had, a concomitant diminution of the HNO₃ should have been observed. Type II PSCs also incorporate HNO₃ vapor, either through the simultaneous cocondensation of HNO₃ and H₂O molecules during ice growth or as NAT particles that serve as condensation nuclei for ice crystals [Toon *et al.*, 1989; Turco *et al.*, 1989]. Thus it seems unlikely that the H₂O mixing ratios could have dropped to relatively low levels while the HNO₃ mixing ratios remained undiminished.

Alternatively, it may be that temperatures fell to sufficiently low values sufficiently quickly that a majority of the background sulfate aerosols froze, precluding the formation of ternary solutions. The presence of at least some frozen sulfate cores would promote the growth of a HNO₃/H₂O solid phase referred to as a type Ic PSC by Tabazadeh and Toon [1996]. Laboratory studies of the deposition of H₂O and HNO₃ vapors onto a solid surface (either glass or SAT) [Hanson, 1992; Marti and Mauersberger, 1993; Iraci *et al.*, 1995] have found evidence for the formation of a metastable water-rich HNO₃-containing solid phase over which the HNO₃ vapor pressure is relatively high (as much as 10–20 times higher than that over NAT). As this HNO₃/H₂O solid phase would also be characterized by a higher HNO₃ vapor pressure than a ternary solution, its formation could explain the initial discrepancy between the ternary model and the MLS data in Plate 18a. In this case, the better agreement between the measured HNO₃ and the NAD model at the end of the observing period could indicate the transfer of vapor from the water-rich solid to NAD as the particles age. A new class of large, water-rich particles has been postulated by Shibata *et al.* [1997] to explain ground-based Arctic lidar measurements that do not conform to the existing type Ia/b PSC definitions. Evidence to support the formation of a water-rich HNO₃/H₂O solid phase has also been reported by Tabazadeh *et al.* [1995] and Tabazadeh and Toon [1996] from an analysis of aircraft observations of PSCs in both the Arctic and the Antarctic. They found the initial depletion of gas-phase HNO₃ due to a type Ic PSC to be only about 2–3 ppbv, which corresponds very closely to the MLS observations shown in Plate 18.

5. Trajectory Calculations

In the PSC formation scenario of Tabazadeh *et al.* [1994a, 1995, 1996] briefly described above, the thermal history of the air mass is critical to the phase of the PSC since it determines the physical state of the preexisting aerosols. These aerosols are liquid in an air mass that either has never been exposed to temperatures low enough for the H₂SO₄ to freeze as SAT (i.e., near the ice frost point [e.g., Middlebrook *et al.*, 1993; Koop *et al.*, 1995]) or has recently experienced

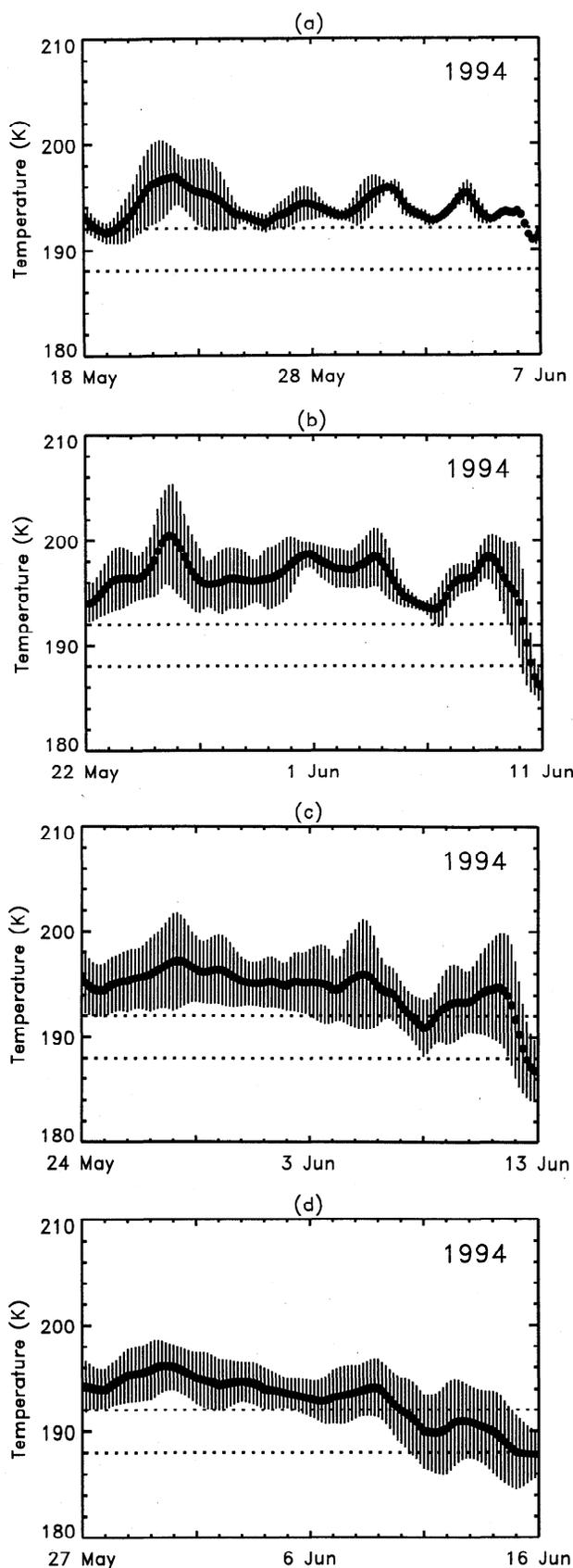
temperatures high enough for SAT to melt (i.e., about 210–220 K [Middlebrook *et al.*, 1993; Zhang *et al.*, 1993b]). On the other hand, the aerosols may freeze through a cooling to temperatures near the ice frost point. In addition, Iraci *et al.* [1994] and Tabazadeh *et al.* [1995] have proposed that after experiencing low temperatures, the H₂SO₄ aerosols may crystallize to SAT upon subsequent warming to about 196–198 K (as observed in the laboratory by Iraci *et al.* [1994]).

To investigate the correlation between synoptic-scale temperature history and PSC composition as inferred from MLS measurements of gas-phase HNO₃, we perform 20-day back trajectory calculations for selected days in each year. As described in more detail by Manney *et al.* [1994], the three-dimensional calculations use horizontal winds and temperatures from the UKMO data assimilation system [Swinbank and O'Neill, 1994]. The trajectory code is formulated in isentropic coordinates and employs a standard fourth-order Runge-Kutta scheme. To account for diabatic effects, a recent version of the middle atmosphere radiation code (MIDRAD) first described by Shine [1987] is used to compute vertical velocities ($d\theta/dt$). Since MLS data do not provide uninterrupted daily coverage, a climatological ozone field is used for the radiation calculations. The impact on the heating rates of using climatological rather than MLS ozone was determined to be negligible by Manney *et al.* [1994]. Although, in general, the large-scale air motion is simulated well by these calculations, Manney *et al.* [1995b] found that they underestimated the magnitude of the diabatic descent in the 1992 Antarctic early winter, based on comparisons with passive tracers measured by CLAES on UARS.

For each trajectory run, parcels are initialized at 465 K in a box defined within the region inside the vortex where HNO₃ mixing ratios are below 4 ppbv; the total number of parcels in the box is chosen to be 100 (10 on a side), 400 (20 on a side), or 900 (30 on a side), depending on the geographic extent of the HNO₃ depletion on that day. The initialization time is 1200 UT, and parcel positions are saved every 4 hours. Temperature histories are then constructed by interpolating UKMO temperatures (available once daily, at 1200 UT) to the time, latitude, longitude, and θ values of these parcels throughout the 20-day run. Overall diabatic effects are found to be essentially similar for runs in all five years, with parcels descending approximately 20 K in the 20 days leading up to the initialization day.

Temperature histories for selected days during the 1994 PSC formation period are shown in Figure 4. Parcels initialized inside the small area of depleted HNO₃ on June 7 (see Plate 1) had experienced a fairly flat synoptic temperature trend (Figure 4a) prior to cooling below 192 K in the last 24 hours. Trajectories run back from the next 6 days exhibit similar features: The parcels experienced large temperature excursions, but they did not undergo significant or sustained cooling until a precipitous drop in temperature (to the ice frost point or below) occurred during the day before the initialization day. This behavior is illustrated in Figure 4b for June 11. The large temperature oscillations in 1994 arose from a combination of the strength and direction of the winds and the shape and location of the low-temperature

areas relative to the vortex. Since a new set of parcels was brought into the cold region each day, ternary solution PSCs continued to form, but the exposure time to low temperatures was too short for the phase transformation to NAD to occur.



This is consistent with the results shown in Plates 5 and 7, which indicate best agreement with the ternary model at this time. By June 13 (Figure 4c), conditions were such that the average temperature of the parcels had briefly dipped below 192 K about 6 days earlier; this may have triggered the formation of some ternary solution PSCs which then converted to NAD, resulting in a mixed cloud as new ternary solutions formed when the temperatures again fell below 192 K on the day before the initialization day. Not until June 16 (Figure 4d) had the parcels initialized within the low-HNO₃ region experienced low temperatures for an extended period of time. Observed HNO₃ for this and subsequent days is most consistent with the NAD model (Plates 5 and 7), suggesting that exposure to low temperatures for at least several days is necessary for significant formation of NAD.

In Figure 5 we show temperature histories for selected days during the early part of the 1993 observing period. Trajectories from the small area of depleted HNO₃ on May 29 (the first day of the observing period; see Plate 12) indicate that these parcels had cooled below 192 K within the last 24 hours (Figure 5a). This pattern of recent cooling below 192 K for parcels originating within the low-HNO₃ region continued through June 1 (Figure 5b). Although some of the aerosols may have undergone crystallization upon warming from ~ 192 K to above 196 K, in general, the temperature histories are such that a majority of the preexisting aerosols were probably in a liquid state, setting the stage for the formation of ternary solution PSCs. By June 2 (Figure 5c), the area of depleted HNO₃ was more extensive (see Plate 12), and the average temperature of the parcels initialized in this region had been continuously at or below 192 K for almost 3 days and had hovered around 192 K for several days before that. Similarly, by June 7 (Figure 5d) the parcels initialized in the low-HNO₃ region had, on average, experienced temperatures at or below 192 K for almost 8 days. Despite the prolonged exposure of these parcels to low temperatures, the model calculations of Plate 14 indicate that better agreement with the data is still achieved at this time using a ternary rather than a NAD composition. As in section 4.3, we attribute this delay in the conversion to NAD to the residual enhancement in aerosol loading from the eruption of Mount Pinatubo.

In both 1992 and 1995, MLS data sampling during early southern winter was limited, and PSC formation was already well underway by the time observations of the southern high latitudes resumed. The extent of HNO₃ loss on the first day of the observing period in these years (June 2 and June 7,

Figure 4. Temperature histories along 20-day back trajectories for the ensemble of parcels located within the region inside the vortex where HNO₃ concentrations are less than 4 ppbv on (a) June 7, 1994, (b) June 11, 1994, (c) June 13, 1994, and (d) June 16, 1994. Dashed lines indicate $T = 188$ K and 192 K. The black circles represent the average temperatures of these parcels at each time step (every 4 hours), and the error bars represent the standard deviations in the averages.

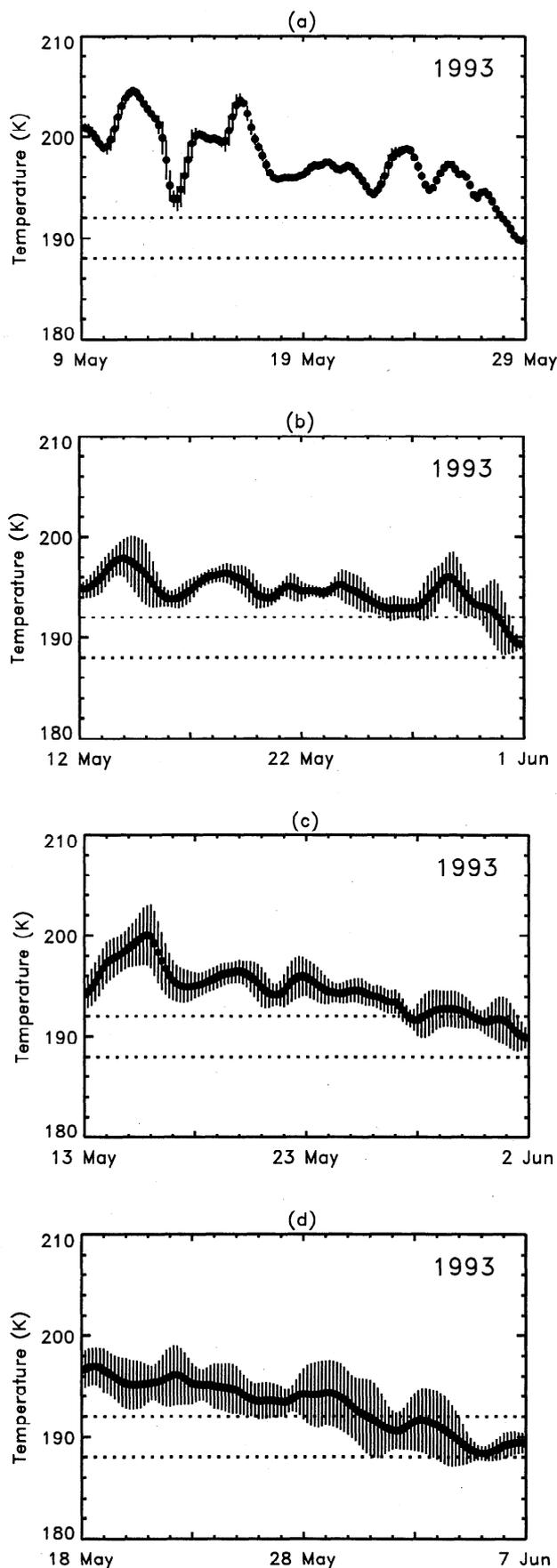


Figure 5. Same as in Figure 4, for (a) May 29, 1993, (b) June 1, 1993, (c) June 2, 1993, and (d) June 7, 1993.

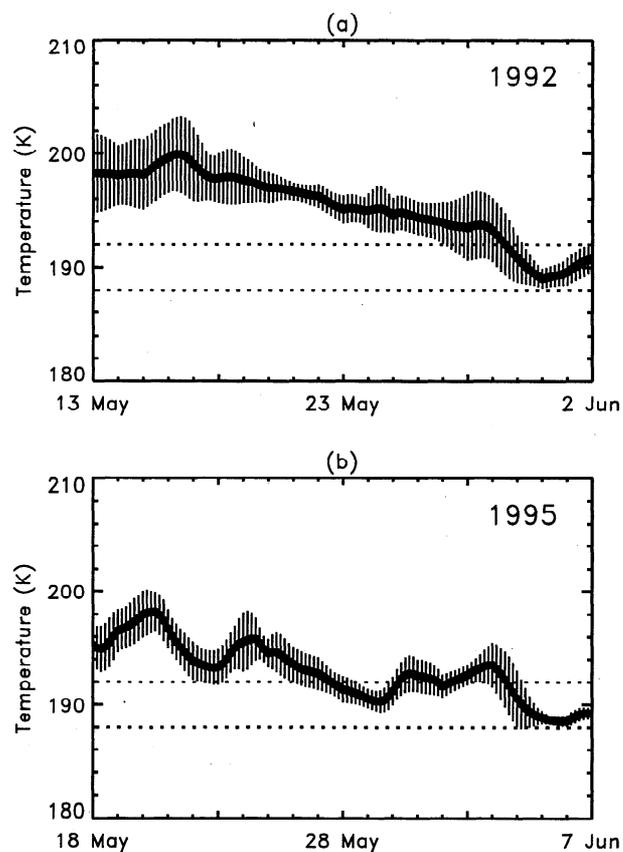


Figure 6. Same as in Figure 4, for (a) June 2, 1992, and (b) June 7, 1995.

respectively) is similar (see Figure 1). In Figure 6a we show the temperature histories of parcels initialized within the low-HNO₃ region on June 2, 1992. In the preceding weeks these parcels experienced a more or less steady temperature decline, and the preexisting aerosols were probably in a liquid state. Just prior to the observation day, the parcels were exposed to temperatures below 192 K for about 3 days, despite a slight warming trend. These conditions are consistent with the formation of ternary solutions. Because 3 days probably would not have allowed sufficient time for NAD nucleation, particularly under the greatly enhanced aerosol loading still present from the Mount Pinatubo eruption (see above), it is likely that the low-HNO₃ observed on June 2 was due to a persistent ternary phase (compare with Plate 10). However, our observations do not exclude a NAD composition. By 1995 the aerosol loading had essentially returned to background levels. The temperature histories of the parcels originating within the low-HNO₃ region on June 7, 1995, are shown in Figure 6b. These parcels had experienced temperatures below 192 K around 8–10 days earlier and had been continuously below 192 K for more than 3 days, allowing time for an initial ternary phase to transform into NAD, as suggested by the observations discussed in section 4.4.

Finally, temperature histories for the 1996 PSC formation period are shown in Figure 7. The average temperature of the parcels initialized inside the small area of depleted HNO₃ on

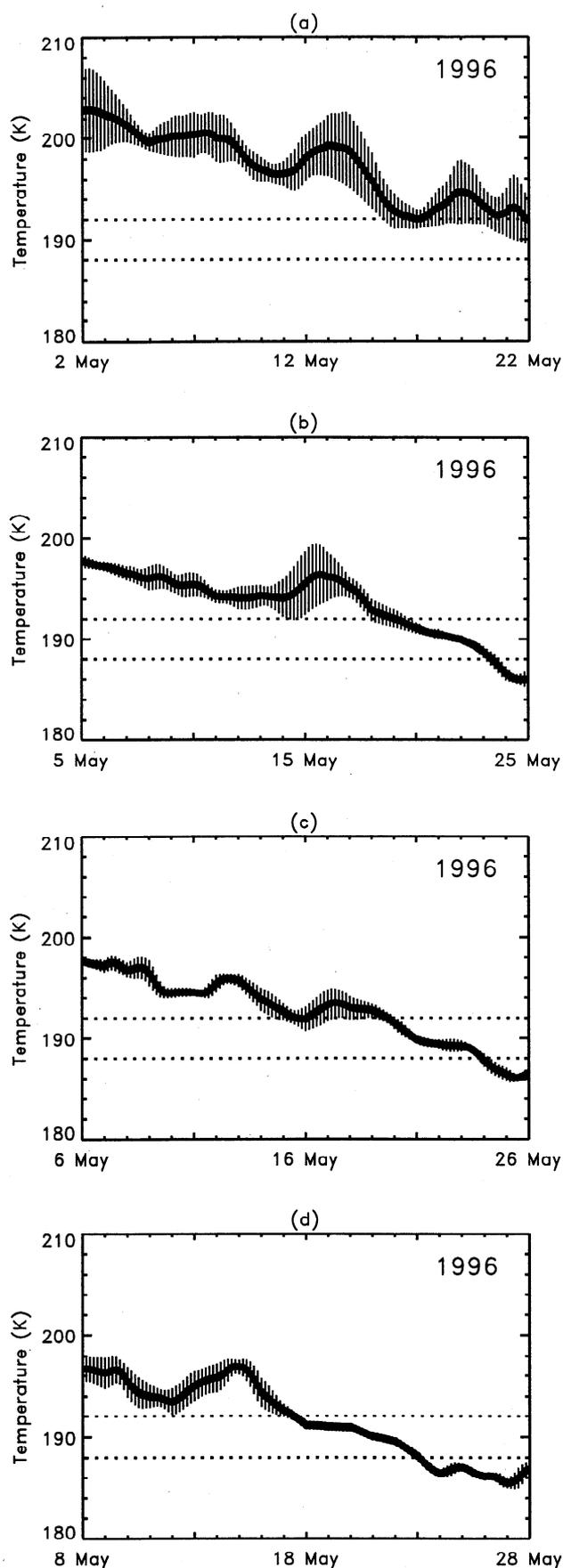


Figure 7. Same as in Figure 4, for (a) May 22, 1996, (b) May 25, 1996, (c) May 26, 1996, and (d) May 28, 1996.

May 22 (see Plate 16) had just dropped to 192 K (Figure 7a); however, the average temperature had been this low on two previous occasions, after which some portion of the parcels encountered temperatures above 196 K. By May 25 and 26 (Figures 7b and 7c), the parcels had been exposed to temperatures below 192 K for more than 6 days and below 188 K for 2 days, and by May 28 (Figure 7d) the parcels had been exposed to temperatures below 192 K for 11 days and below 188 K for 5 days; yet in all these cases the MLS data indicate higher gas-phase HNO₃ values than predicted by any of the PSC composition models (see Plate 18). The parcels originating within the low-HNO₃ region on these days experienced larger sustained synoptic cooling rates and spent more time at temperatures near the ice frost point than in previous years. Hence a larger fraction of the sulfate aerosols may have frozen in the early winter of 1996, inhibiting the growth of ternary solutions. The presence of frozen sulfate cores would lead to the emergence of type Ic PSCs, whose relatively high HNO₃ vapor pressure provides the best match to the pattern of depletion seen in the 1996 MLS HNO₃ data.

6. Summary and Implications

We have examined MLS HNO₃ measurements obtained at the beginning of five southern hemisphere winters: 1992–1996. The observed evolution of the gas-phase HNO₃ was compared against that predicted by NAT [Hanson and Mauersberger, 1988], NAD [Worsnop *et al.*, 1993], and liquid ternary solution [Tabazadeh *et al.*, 1994b] models and correlated with temperature histories from three-dimensional back trajectory calculations to infer the composition of the PSCs that formed in early winter each year.

MLS observations of the changes in gas-phase HNO₃ suggest that the first PSCs to form are composed of either liquid ternary solutions or a metastable water-rich solid phase, depending on the physical state of the background sulfate aerosols. This is consistent with the PSC formation scenario of Tabazadeh *et al.* [1994a, 1995, 1996]. According to this scenario, if the preexisting aerosols are liquid, then ternary solutions form as the temperature drops below about 192 K, a few degrees below the NAT condensation point. After prolonged exposure to low temperatures, NAT or NAD particles begin to nucleate, and a mixed cloud results while a transfer of HNO₃ vapor from the ternary solution droplets to the more stable crystalline particles takes place, with eventual complete conversion. This paradigm can explain the MLS 1994 observations, which show better agreement with the ternary solution model in the initial stages of PSC formation but which are matched best by the NAD model as time progresses and the air masses have been exposed to low temperatures for several days. Although the 1995 MLS data are too sparse to allow a definitive conclusion, they also appear to be consistent with this PSC formation scenario.

The basic PSC formation process outlined above also applies to MLS measurements from 1993, when the stratospheric particle optical depth was still elevated by roughly a factor of 10 following the eruption of Mount Pinatubo

[Russell *et al.*, 1996]. However, the 1993 MLS data indicate that HNO₃ uptake into ternary solutions occurs at higher temperatures under enhanced aerosol conditions, as predicted by Carslaw *et al.* [1994] and Tabazadeh *et al.* [1994b]. In addition, the data suggest that the conversion to NAD is delayed under volcanic conditions, since even after the air masses have experienced temperatures below 192 K for almost 8 days the ternary model still provides a better match to the data than the NAD model. This is probably related to the significant differences in composition between ternary solutions formed under volcanic conditions and those formed under background aerosol conditions [Tabazadeh *et al.*, 1994b]. These compositional differences during volcanically perturbed periods may result in changes in the NAD saturation ratio that hinder the formation of crystalline PSCs. After exposure to low temperatures for about 2 weeks, however, the 1993 MLS data are matched best by the NAD model. The MLS data for 1992, when the stratospheric particle optical depth was enhanced by roughly a factor of 20–30 over pre-Pinatubo levels [Russell *et al.*, 1996], may also indicate a persistent ternary phase, but again, the data are too sparse to conclusively distinguish between NAD and ternary solution compositions.

MLS measurements from 1996 indicate less HNO₃ depletion than predicted by any of the PSC composition models. During the 1996 early southern winter the lower stratosphere was substantially colder than average, and the air masses participating in the initial HNO₃ depletion experienced larger sustained synoptic cooling rates and spent more time at temperatures near the ice frost point than in the previous years. Hence a majority of the sulfate aerosols may have frozen in the 1996 early winter, inhibiting the growth of ternary solutions. According to the PSC formation mechanism of Tabazadeh *et al.* [1994a, 1995, 1996], background aerosols in a frozen state promote the initial formation of a metastable water-rich HNO₃-containing solid phase (type Ic), over which the HNO₃ vapor pressure is relatively high. The 1996 MLS HNO₃ data are consistent with the magnitude of the gas-phase HNO₃ depletion attributed to type Ic PSCs by Tabazadeh and Toon [1996]. Better agreement between the measured HNO₃ and the NAD model at the end of the observing period may indicate the transfer of vapor from the water-rich solid to NAD as the particles age. Some type Ic PSCs may also have formed in the previous years studied, but their presence would have been masked by the much greater uptake of HNO₃ into ternary solutions if most of the background aerosols were liquid.

In general, we have found a strong correspondence between the area of gas-phase HNO₃ loss and the area of temperatures below 192 K but only a weak correspondence between the area of gas-phase HNO₃ loss and the area of temperatures below 195 K, the value commonly assumed as the threshold for PSC formation. In fact, for the 3 years in which MLS made measurements at or near the time of initial PSC formation (1993, 1994, and 1996), large NAT supersaturations persisted for days to weeks before significant depletion of gas-phase HNO₃ occurred. Although the temperatures were low enough to maintain NAT PSCs, the MLS data in-

dicate that they were not forming, at least not over spatial scales comparable to or larger than the $\sim 400 \times 100 \times 5$ km MLS field of view.

Although our results are strongly suggestive, the limited precision and spatial resolution of the data and the necessity of estimating unmeasured quantities preclude conclusive identification of PSC compositions. With this caveat in mind, however, it appears that the ternary solution and NAD models provide good agreement with the MLS HNO₃ measurements at different stages of the PSC formation process, while the NAT model consistently overestimates the observed removal of gas-phase HNO₃. This has important implications for theoretical studies of ozone depletion. Many stratospheric models employ PSC parameterizations whereby it is assumed that NAT particles form and heterogeneous processing begins immediately after temperatures drop below 195 K. Since most such models have a horizontal resolution similar to that of the MLS data, our results suggest that these models may not calculate ozone loss accurately.

The majority of the MLS HNO₃ measurements obtained during the early southern winter periods indicate that exposure to low temperatures for at least several days (longer under conditions of enhanced aerosol loading following volcanic eruptions) is necessary for a significant degree of crystalline PSC formation. This suggests the existence of substantial free-energy barriers to nucleation for the crystalline forms. That the MLS data agree better with NAD for aged PSCs is supportive of the suggestion of a lower barrier to nucleation for NAD than for NAT [Worsnop *et al.*, 1993; Düsselkamp *et al.*, 1996]. In any case, large barriers to nucleation would allow activation and growth of only a small fraction of available condensation nuclei (regardless of cooling rates) and so would result in formation of a small number of relatively large HNO₃-containing particles that could fall appreciable distances, leading to denitrification [Salawitch *et al.*, 1989].

In situ [e.g., Fahey *et al.*, 1990] and satellite [Roche *et al.*, 1994; Santee *et al.*, 1995] observations have shown that both the temporal and the spatial scale of denitrification are substantially greater in the Antarctic than in the Arctic, consistent with the generally warmer and weaker lower stratospheric vortex [e.g., Andrews, 1989] and fewer, shorter-duration PSC occurrences [Poole and Pitts, 1994] in the northern hemisphere. However, enhanced ClO abundances, arising through reactions on the surfaces of PSC particles, comparable to those over Antarctica are routinely observed throughout the Arctic polar vortex [e.g., Waters *et al.*, 1993; Santee *et al.*, 1995]. Interhemispheric differences in the severity of ozone loss have been ascribed to the moderating influence of high levels of HNO₃ on active chlorine during the latter stages of Arctic winter [e.g., Brune *et al.*, 1991; Santee *et al.*, 1995]. The MLS HNO₃ data presented here suggest that small liquid droplets form quickly at temperatures below ~ 192 K, initiating chlorine activation, but air masses must be exposed to low temperatures for several days or longer to facilitate growth of the large crystalline particles that undergo significant gravitational settling. Since in

the Arctic the low temperatures are usually confined to relatively small regions that are not concentric with the vortex, individual air masses rarely experience temperatures below 192 K for extended periods of time. Thus the much smaller degree of denitrification in the Arctic is probably related not only to the lack of temperatures below the water ice frost point (~188 K) but also to the lack of temperatures persistently below ~192 K.

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- L. Froidevaux, G. L. Manney, W. G. Read, R. J. Salawitch, M. L. Santée (corresponding author), and J. W. Waters, Jet Propulsion Laboratory, Mail Stop 183–701, 4800 Oak Grove Drive, Pasadena, CA 91109. (e-mail: mls@mls.jpl.nasa.gov)
- A. Tabazadeh, NASA Ames Research Center, Moffett Field, CA 94035. (e-mail: taba@sky.arc.nasa.gov)

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